

Origin of Mass and Scattering Amplitudes: from Higgs to Pauli, Kaluza-Klein and Chern-Simons

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Wuhan, December 15, 2025

References

➤ Based on series of works:

You, Song, HJH, Han, PRD(2025), 2412.16033

Han, HJH, You, Song, PRD-Lett (2025), 2412.21045

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Y. Li, Y.F. Hang, HJH, S. He, JHEP (2022), 2111.12042 [hep-th].

H.X. Liu, Z.X. Yi, HJH, [arXiv:2512.10870 \[hep-th\]](#) (66 pp).

Y.F. Hang, HJH, C. Shen, JHEP (2022), 2110.05399 [hep-th].

Y.F. Hang, HJH, C. Shen, Research (2023), 2406.13671 [hep-th].

➤ Early works:

Chivukula, Dicus, HJH, PLB (2002) [hep-ph/0111016]

Chivukula, HJH, PLB (2002) [hep-ph/0201164]

H.J. He, IJMP (2005), APS-2004 [hep-ph/0412113]

Chivukula, Dicus, HJH, Nandi, PLB (2003) [hep-ph /0302263]

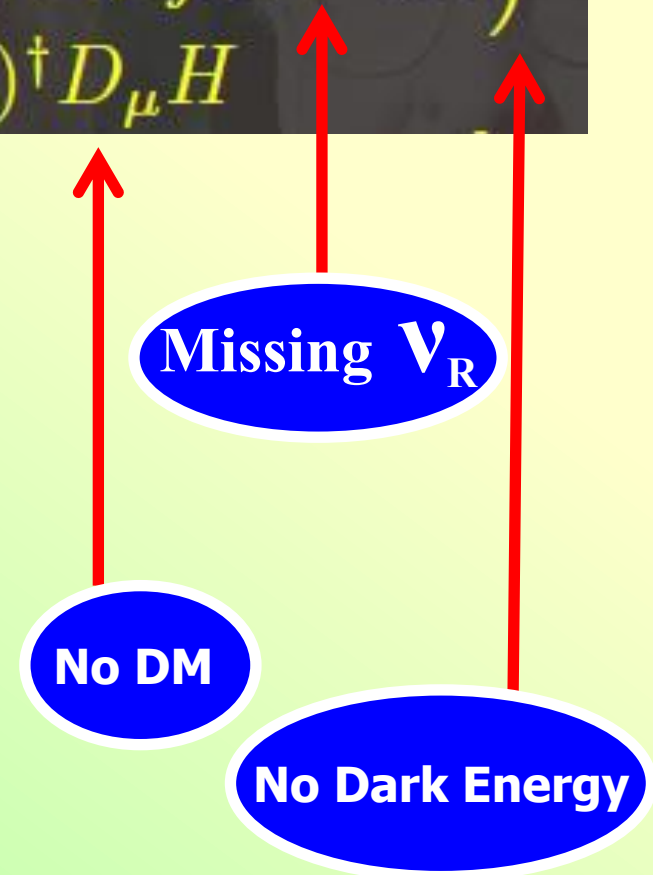
Making of the Standard Model (123)

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4g'^4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4g^2} W_{\mu\nu}^a W^{\mu\nu a} - \frac{1}{4g_s^2} G_{\mu\nu}^a G^{\mu\nu a} \\ & + \bar{Q}_i i \not{D} Q_i + \bar{u}_i i \not{D} u_i + \bar{d}_i i \not{D} d_i + \bar{L}_i i \not{D} L_i + \bar{\ell}_i i \not{D} \ell_i \\ & + \left(Y_u^{ij} \bar{Q}_i u_j \tilde{H} + Y_d^{ij} \bar{Q}_i d_j H + Y_l^{ij} \bar{L}_i \ell_j H + c.c. \right) \\ & - \lambda (H^\dagger H)^2 + \lambda v^2 H^\dagger H - (D^\mu H)^\dagger D_\mu H\end{aligned}$$

Is **SM** Structure complete? → **not really !!**

→ What is really **Missing** in the SM ???

- 1) ν_R Within SM Structure **not yet found !**
- 2) **No Dark Matter** (? X)
- 3) **No Dark Energy** (? ✓)



Making of the Standard Model (123)

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4g'^4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4g^2} W_{\mu\nu}^a W^{\mu\nu a} - \frac{1}{4g_s^2} G_{\mu\nu}^a G^{\mu\nu a} \\ & + \bar{Q}_i i \not{D} Q_i + \bar{u}_i i \not{D} u_i + \bar{d}_i i \not{D} d_i + \bar{L}_i i \not{D} L_i + \bar{\ell}_i i \not{D} \ell_i \\ & + \left(Y_u^{ij} \bar{Q}_i u_j \tilde{H} + Y_d^{ij} \bar{Q}_i d_j H + Y_l^{ij} \bar{L}_i \ell_j H + c.c. \right) \\ & - \lambda (H^\dagger H)^2 + \lambda v^2 H^\dagger H + (D^\mu H)^\dagger D_\mu H\end{aligned}$$

SM Structure is NOT yet complete

→ What is *Missing* in the SM ??

→→ **No full understanding on Quantum Gravity**
at **both the Largest and Smallest Scales !!**

→→ UV-IR correspondence ? ...

No Gravity!

Missing ν_R

Origin of Neutrino Mass Generation and How to Probe it ?

**You, Song, HJH, Han, PRD(2025), 2412.16033
Han, HJH, You, Song, PRD-Lett (2025), 2412.21045**

Wolfgang Pauli: Father of the Neutrinos



W. Pauli (1900-1958)

- 1914: James Chadwick discovered that β decay has continuous electron energy spectrum and could not explain it with 2-body final state.
- In 1929, Niels Bohr attempted to explain this by giving up energy conservation (X).
- 1930: Pauli postulated that the final state is 3-body, containing a new massless particle with electric charge 0 & spin-1/2. This naturally explains the continuous E-spectrum + conserves energy. (✓)
- 1932: Fermi named this particle as “Neutrino”.
1933: Fermi wrote down 4-Fermion Interaction to describe β decays in weak interaction.



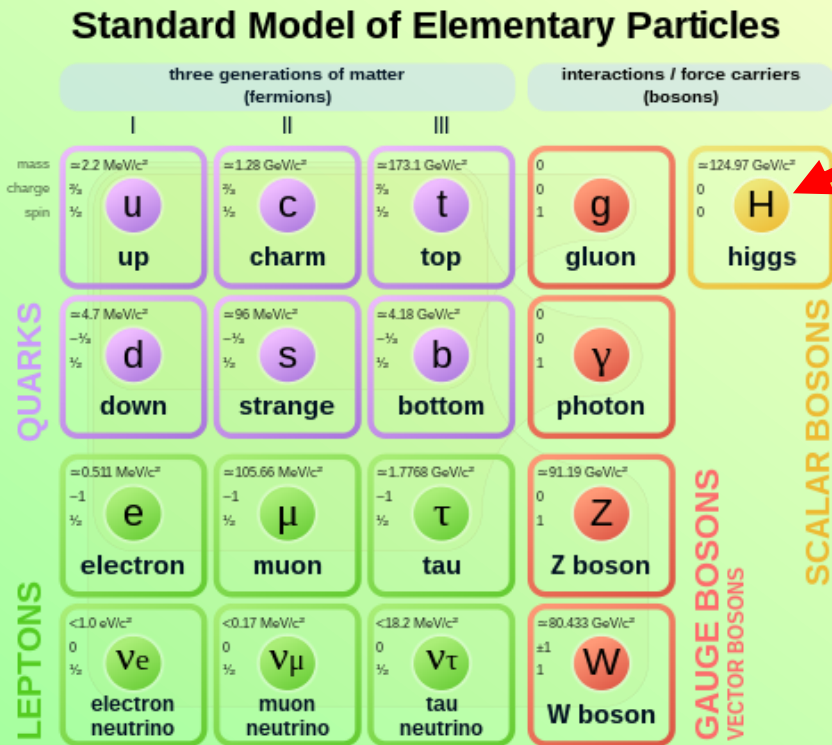
E. Fermi (1901-1954)

Last Missing Piece of the SM

- For conventional SM setup before 1998, neutrinos were assumed **for simplicity** to be **massless** and have only left-handed components because the SM is structured to have **all the right-handed fermions be weak singlets** in each fermion family, where the **right-handed neutrinos (ν_R) are pure gauge singlets** and **their absence does not affect the gauge anomaly cancellation of the SM**.
- Weinberg realized that without ν_R , the left-handed neutrinos can acquire small **Majorana masses** from a **gauge-invariant dimension-5 operator (LLHH)** that is suppressed by a large UV cutoff scale $\Lambda_\nu \sim v^2/m_\nu$, **far beyond the weak scale**.
- However, this dimension-5 operator is **nonrenormalizable** and its **Minimal UV Completion** is given by the **canonical seesaw** with $\Lambda_\nu = M_R$ after adding back ν_R for each fermion family.
- The existence of ν_R is **predicted by the SM structure** and provides the **minimal UV completion for dimension-5 Weinberg operator** through **canonical Seesaw Mechanism**, naturally generating light neutrino masses.
- Yet, ν_R points to a **brand-new seesaw scale** $\Lambda_\nu \sim v^2/m_\nu \sim 10^{14}\text{GeV}$ that is **far beyond SM**.
- $\rightarrow\rightarrow$ It is extremely important to probe ν_R as the **Last Missing Piece of the SM** and test **Neutrino Mass-Generation** via **canonical Seesaw Mechanism**.
- $\rightarrow\rightarrow$ SM prediction ν_R could be wrong by **EXP**, but again **its chance of success is high !**

Structures of 2 SM

SM of Particle Physics + SM of Cosmology (粒子物理标准模型 + 宇宙学标准模型)



Higgs **$h(125\text{GeV})$** (2012) is actually **NOT** the Last Building Block of SM .

But, **ν_R** is the **Last Missing Piece of SM !!**

GR - Einstein-Hilbert Action:

$$S = \int d^4x \sqrt{-g} [\kappa^{-2}(R + \Lambda) + \mathcal{L}_m]$$

SM of Particle Physics + SM of Cosmology

粒子物理标准模型 + 宇宙学标准模型

Standard Model of Elementary Particles

three generations of matter (fermions)			interactions / force carriers (bosons)	
I	II	III		
mass charge spin				
$\approx 2.2 \text{ MeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	$\approx 1.28 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	$\approx 173.1 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	0 0 1	$\approx 124.97 \text{ GeV}/c^2$ 0 0
u up	c charm	t top	g gluon	H higgs
$\approx 4.7 \text{ MeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$	$\approx 96 \text{ MeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$	$\approx 4.18 \text{ GeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$	0 0 1	
d down	s strange	b bottom	γ photon	
$\approx 0.511 \text{ MeV}/c^2$ -1 $\frac{1}{2}$	$\approx 105.66 \text{ MeV}/c^2$ -1 $\frac{1}{2}$	$\approx 1.7768 \text{ GeV}/c^2$ -1 $\frac{1}{2}$	0 0 1	
e electron	μ muon	τ tau	Z Z boson	
$< 1.0 \text{ eV}/c^2$ 0 0	$< 0.17 \text{ MeV}/c^2$ 0 $\frac{1}{2}$	$< 18.2 \text{ MeV}/c^2$ 0 $\frac{1}{2}$	$\approx 80.433 \text{ GeV}/c^2$ ± 1 1	
ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

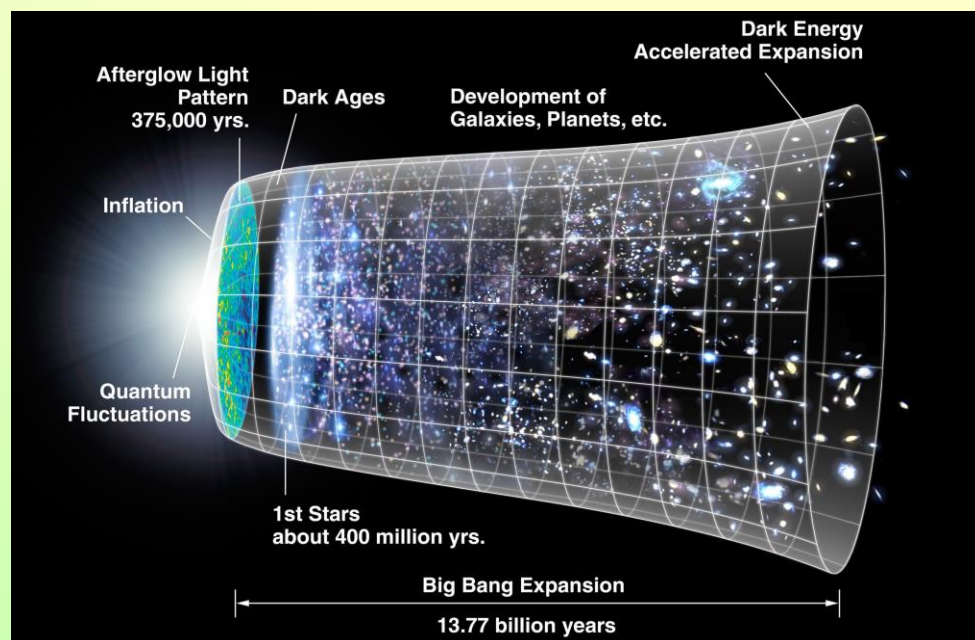
QUARKS

LEPTONS

GAUGE BOSONS
VECTOR BOSONS

SCALAR BOSONS

Cosmological SM: Λ CDM + Inflation



物质的起源与演化

Missing ν_R

No DM??

No Dark Energy??

Cosmological Bounds on **light ν Mass Sum**:

Table 26.2: Summary of $\sum m_\nu$ constraints.

PDG

	Model	95% CL (eV)	Ref.
CMB alone			
Pl18[TT+lowE]	$\Lambda\text{CDM}+\sum m_\nu$	< 0.54	[24]
Pl18[TT,TE,EE+lowE]	$\Lambda\text{CDM}+\sum m_\nu$	< 0.26	[24]
CMB + probes of background evolution			
Pl18[TT,TE,EE+lowE] + BAO	$\Lambda\text{CDM}+\sum m_\nu$	< 0.13	[49]
Pl18[TT,TE,EE+lowE] + BAO	$\Lambda\text{CDM}+\sum m_\nu+5 \text{ params.}$	< 0.515	[25]
CMB + LSS			
Pl18[TT+lowE+lensing]	$\Lambda\text{CDM}+\sum m_\nu$	< 0.44	[24]
Pl18[TT,TE,EE+lowE+lensing]	$\Lambda\text{CDM}+\sum m_\nu$	< 0.24	[24]
Pl18[TT,TE,EE+lowE]+ ACT[lensing]	$\Lambda\text{CDM}+\sum m_\nu$	< 0.12	[50]
CMB + probes of background evolution + LSS			
Pl18[TT,TE,EE+lowE] + BAO + RSD	$\Lambda\text{CDM}+\sum m_\nu$	< 0.10	[49]
Pl18[TT,TE,EE+lowE+lensing] + BAO + RSD + Shape	$\Lambda\text{CDM}+\sum m_\nu$	< 0.082	[51]
Pl18[TT+lowE+lensing] + BAO + Lyman- α	$\Lambda\text{CDM}+\sum m_\nu$	< 0.087	[52]
Pl18[TT,TE,EE+lowE] + BAO + RSD + SN + DES-Y1	$\Lambda\text{CDM}+\sum m_\nu$	< 0.12	[49]
Pl18[TT,TE,EE+lowE] + BAO + RSD + SN + DES-Y3	$\Lambda\text{CDM}+\sum m_\nu$	< 0.13	[53]

→ Finally, the **heaviest Light Neutrino Mass** should be around **0.05eV (IO) — 0.06eV (NO)**.

Upper Bounds on **Largest Light ν Mass**

$$\begin{aligned}\text{NO: } \sum m_\nu &= m_3 + \sqrt{m_3^2 - \Delta m_{31}^2} + \sqrt{m_3^2 - \Delta m_{31}^2 + \Delta m_{21}^2} \\ \text{IO: } \sum m_\nu &= m_2 + \sqrt{m_2^2 - \Delta m_{21}^2} + \sqrt{m_2^2 - |\Delta m_{32}|^2},\end{aligned}$$

➤ **Largest Light Neutrino Mass:**

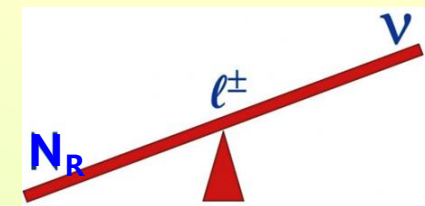
→→ **Normal Ordering (NO): 0.06 eV**

→→ **Inverted Ordering (IO): 0.05 eV**

Beyond SM: **Scale of Neutrino Mass Generation**

- Weinberg Operator: $1/\Lambda_\nu$ (LLHH)
→→ **Model-independent** formulation of ν Mass Generation !
- **Canonical Seesaw is the Minimal UV Completion.**
- **Seesaw Scale:** $\Lambda_\nu = M_R = m_D^2/m_\nu \sim 10^{14} \text{ GeV}$
→→ **Seesaw** Scale is a **brand-new Scale** beyond SM !

Neutrino Seesaw



➔ **Great Challenge:**

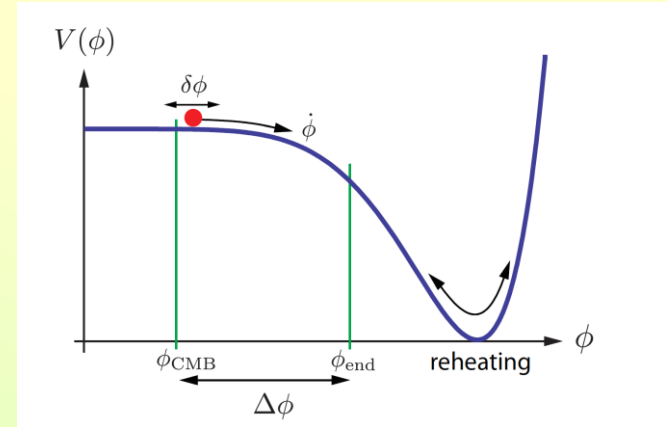
How to test the **High Scale Seesaw** ??

Minimal Model

You, Song, HJH, Han, PRD(2025), 2412.16033
Han, HJH, You, Song, PRDLett (2025), 2412.21045

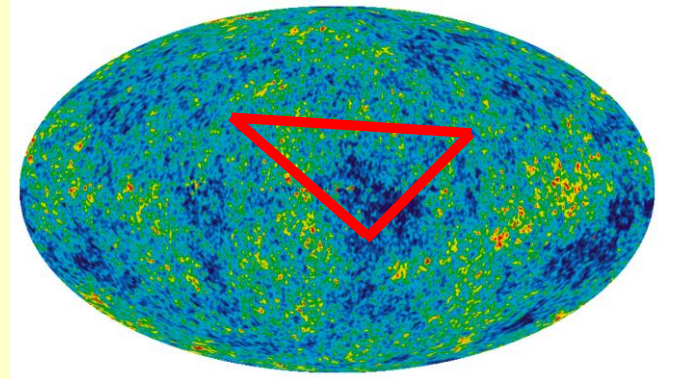
Inflation + Seesaw occur around Same Scale $\sim 10^{14}$ GeV !

$$\Delta\mathcal{L} = \sqrt{-g} \left[-\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) + \bar{N}_R i \not{\partial} N_R + \frac{1}{\Lambda} \partial_\mu \phi \bar{N}_R \gamma^\mu \gamma^5 N_R \right. \\ \left. + \left(-\frac{1}{2} M \bar{N}_R^c N_R - y_\nu \bar{\ell}_L \tilde{H} N_R + \text{H.c.} \right) \right],$$



- $V(\phi)$ is the potential for inflation, is unknown but dominated by the mass term after inflation.
- Derivative coupling to keep the flatness of the inflaton potential (shift-symmetry).
- $\Lambda > 60 H_{\text{inf}}$ to keep perturbative unitarity.
- Inflaton coupling with SM fermions does not affect the analysis.
- After inflation, inflaton oscillates at the bottom of the potential until it decays into heavy N_R ($M_\phi > 2 M_R$). The heavy neutrinos quickly decay into SM particles and reheat the universe.

Non-Gaussianity



Non-Gaussianity is **sensitive** to **New Physics** !

- **New physics could induce large non-Gaussianity:** multi-field inflation models, modulated reheating, curvaton scenario
- **Current limit from Planck on local type $f_{\text{NL}} \sim \mathcal{O}(10)$, future CMB observations, LiteBIRD $\mathcal{O}(1)$, large scale structure observations DESI $\mathcal{O}(1)$, 21cm tomography $\mathcal{O}(0.01-0.1)$**
- **Non-Gaussianity can provide information to the new particle mass, spin, interactions:** **Cosmological Collider Signals**

Local type non-Gaussianity

- The local type non-Gaussianity which is defined by Bardeen Potential:

$$\Phi \equiv \frac{3}{5}\zeta$$

$$\langle \Phi_{\mathbf{k}_1} \Phi_{\mathbf{k}_2} \Phi_{\mathbf{k}_3} \rangle'_{\text{local}} = 2A^2 f_{\text{NL}}^{\text{local}} \left\{ \frac{1}{k_1^3 k_2^3} + \frac{1}{k_2^3 k_3^3} + \frac{1}{k_3^3 k_1^3} \right\}$$

- In the limit $k_1 \sim k_2 \gg k_3$, we find

$$f_{\text{NL}}^{\text{local}} \sim -\frac{10}{3} \frac{z_1^3 H^3}{(2\pi)^4 \mathcal{P}_\zeta^2} \cdot \left(\frac{\lambda \bar{h}}{2H} N_e - \frac{H \cdot z_2}{4z_1} \right)$$

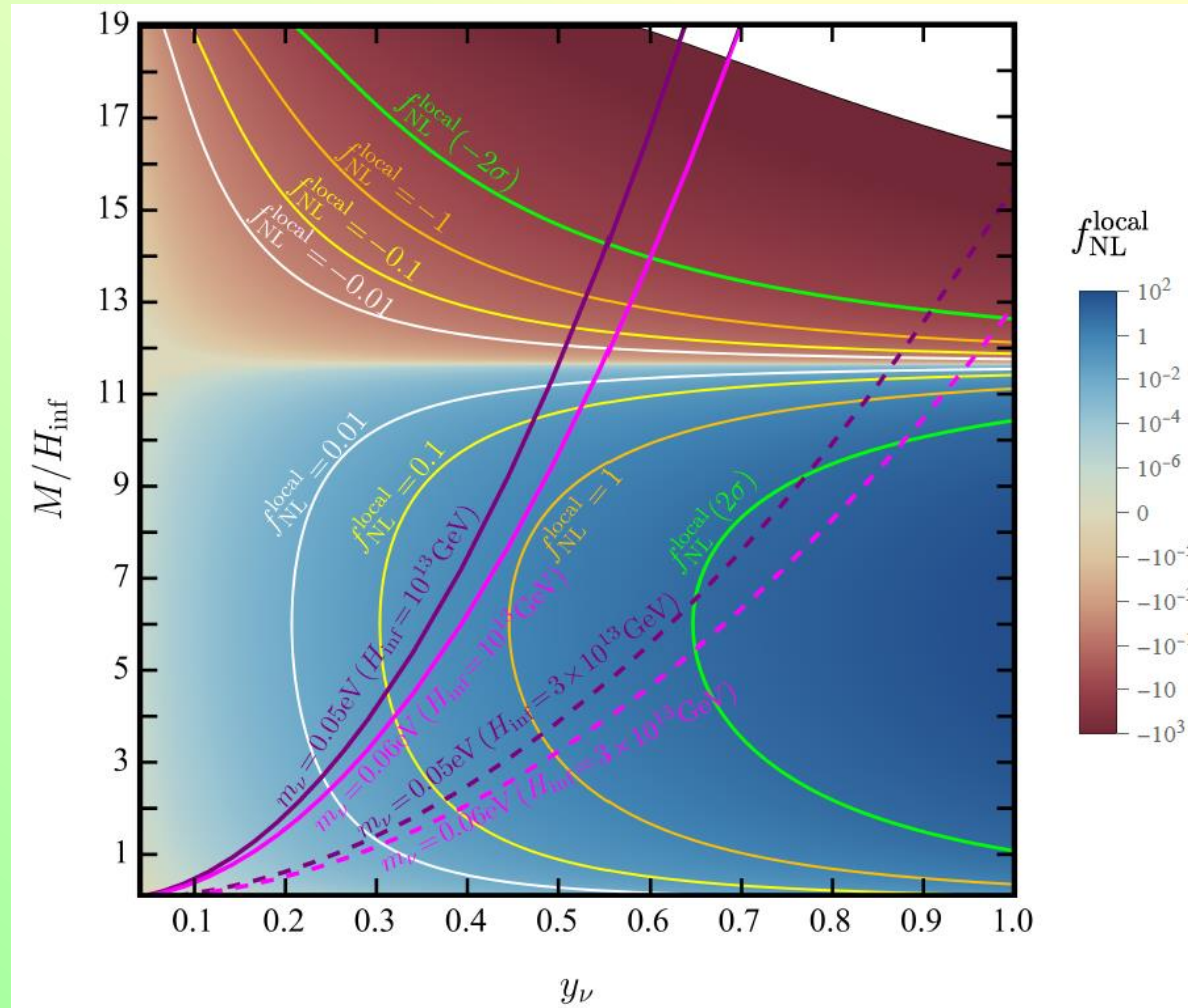
$$f_{\text{NL}}^{\text{local}} = -0.9 \pm 5.1 \quad (68\% \text{ C.L., Planck 2018})$$

- Benchmark for this analysis:

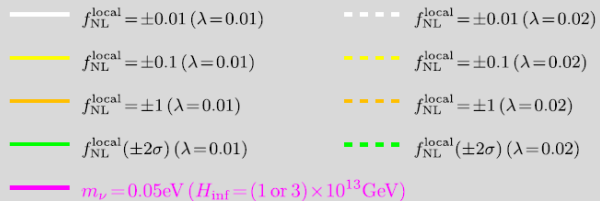
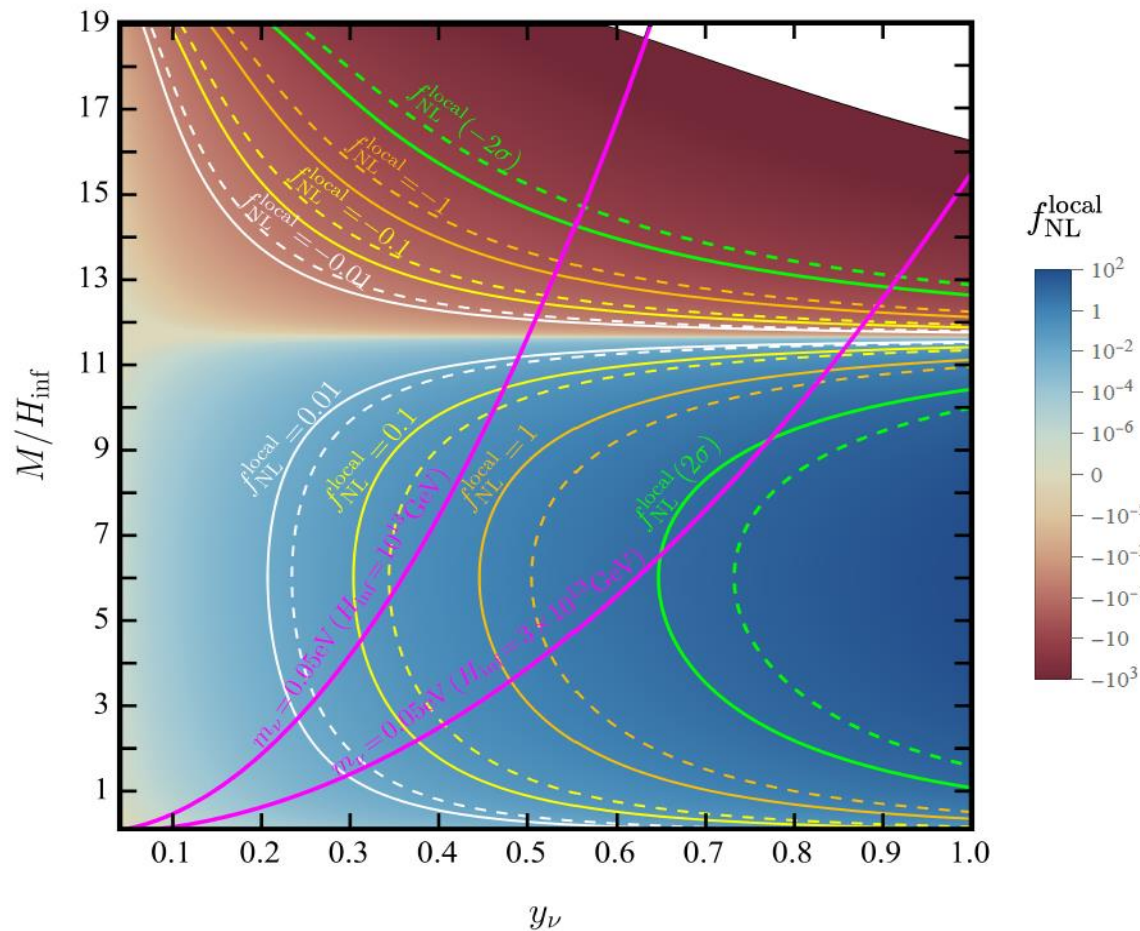
Parameters	\mathcal{P}_ζ	N_e	H_{inf}	m_ϕ	Λ	λ
Values	2.1×10^{-9}	60	$(1, 3) \times 10^{13} \text{ GeV}$	$40 H_{\text{inf}}$	$60 H_{\text{inf}}$	0.01

non-Gaussianity Probes Seesaw:

M_R vs Yukawa y_ν
and ν_R vs ν_L Mass Scales



NG: Seesaw & Higgs Self-Coupling



Summary-1

- We propose a Minimal Model **incorporating inflation and seesaw** through $\phi - \nu_R$ interaction.
It provides a **New Realization** of **Higgs-modulated reheating**.
- It gives a **direct probe** of **Seesaw** for **Neutrino Mass Generation !**
It links **Non-Gaussianity constraint** to **Low Energy ν -Exps** on **light ν masses** and to **LHC measurements** on **Higgs self-coupling**, via RG running.
- **Seesaw generated non-Gaussianity** could be probed in near future CMB or large-scale structure observations
- It also provides a framework of cosmological Higgs collider (**particles coupling to Higgs boson could be detected**)

Geometric Higgs Mechanism for KK Graviton Mass Generation

Hang and HJH, PRD(2022), 2106.04568

Research (2022), 2207.11214

Hang, Zhao, HJH, Qiu, (2024), 2406.12713

Challenges of Mass Generations in 4d

Nothing forbids **Spin-0 Higgs Boson** gets a huge mass from quantum corrections!
Quark/Lepton mass-scales are protected by P-violation and set by Higgs VEV through Yukawa coupling.

Fields	$m = 0$		$m \neq 0$
→ H^0 ($s=0$)	1 ($\lambda = 0$)	=	1 ($\lambda = 0$)
A^μ ($s=1$)	2 ($\lambda = \pm 1$)	\neq	3 ($\lambda = 0, \pm 1$)
$h^{\mu\nu}$ ($s=2$)	2 ($\lambda = \pm 2$)	\neq	5 ($\lambda = 0, \pm 1, \pm 2$)
→ (ψ_L, ψ_R) ($s = \frac{1}{2}$)	1 + 1 ($\lambda = \frac{1}{2}, -\frac{1}{2}$)	=	2 ($\lambda = \pm \frac{1}{2}$)

Higgs Mechanism and Beyond

➤ “Higgs” Mechanism can be

1). **Conventional**, 2). **Geometric**, 3). **Topological**

-- **Conventional Higgs Mechanism:**

Mass-Generation by SSB by Vacuum Expectation Value of Higgs Boson.

-- **Geometric “Higgs” Mechanism:**

Mass-Generation SSB by Compactification of Extra Dimension.

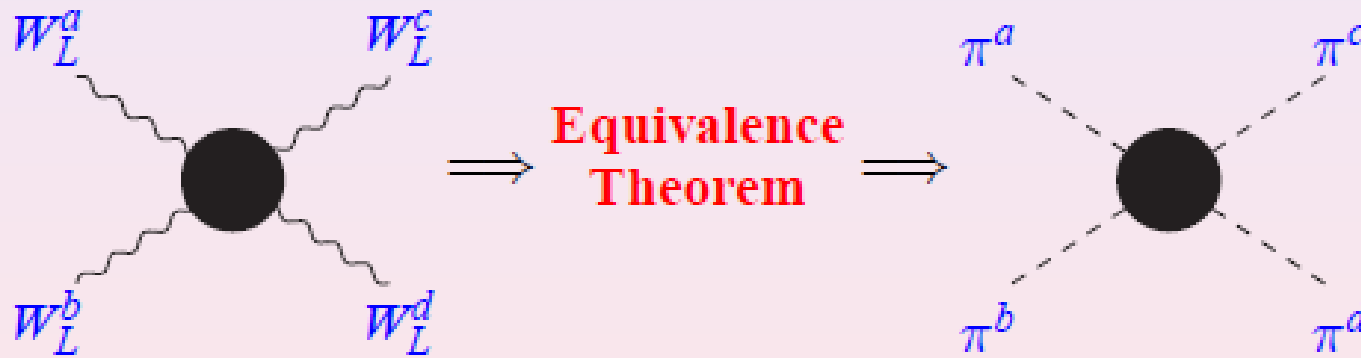
-- **Topological “Higgs” Mechanism:**

No SSB, construction of **Gauge-Invariant Mass Term** with topology,
Dynamical Conversion of Physical Degrees of Freedom.

SM: **ET** as **Bridge** between **Exp** and **EWSB**

No Lose Theorem for LHC Discovery

— **Equivalence Theorem** —
the Bridge between Experiments and EWSB



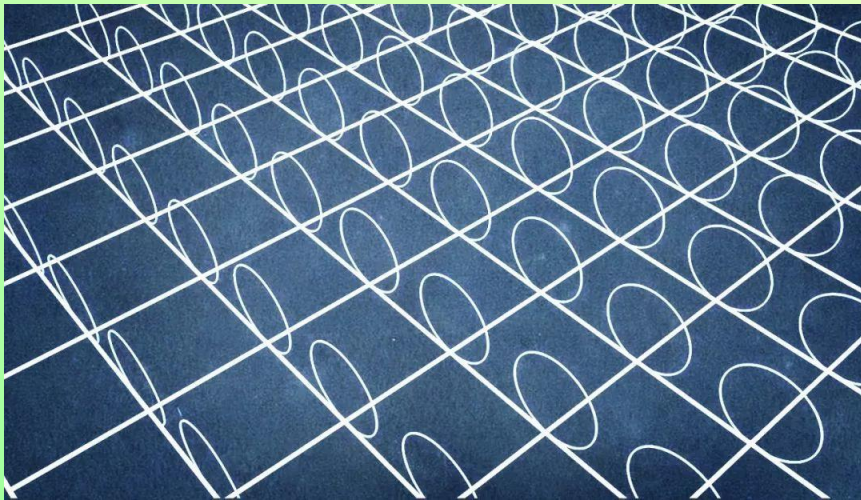
Energy Cancellations: $E^4 \rightarrow E^0$

Unitarity Bound $\rightarrow m_H < 800\text{GeV}$ or $\Lambda < 1.2\text{TeV}$

For a review: H.-J. He et al, arXiv:hep-ph/9704276.

5d: Kaluza-Klein Compactification

- We are apparently living in $(3+1)d$ spacetime.
- But our Universe could have **Extra Dimensions** beyond $d = 4$, except that **Extra Dimensions are all curled-up !**
- Simplest case is a 5th dimension curled up on a circle S^1 , called **Kaluza-Klein (KK) Compactification** (in 1920s)



Geometric “Higgs” Mechanism in 5d

➤ Under compactification of flat 5d, how do **KK Masses** arise? ---

▶ Consider 5d Massless Gauge Fields $\hat{A}^{aM}(x^\mu, x^5)$,

$$\begin{cases} 0 = P^2 = p_\mu p^\mu + p_5 p^5 = p^2 - p_5^2 \\ p^5 = \frac{n}{R}, \quad (n = 0, 1, 2, \dots), \end{cases}$$

(After 5d compactification)

where p^5 is quantized due to BCs, and $p^2 = p_\mu p^\mu$ is the 4-momentum-squared in 4d.

▶ Hence, in 4d we can see the “**KK Tower**”, as a unique consequence of **5d Compactification**,

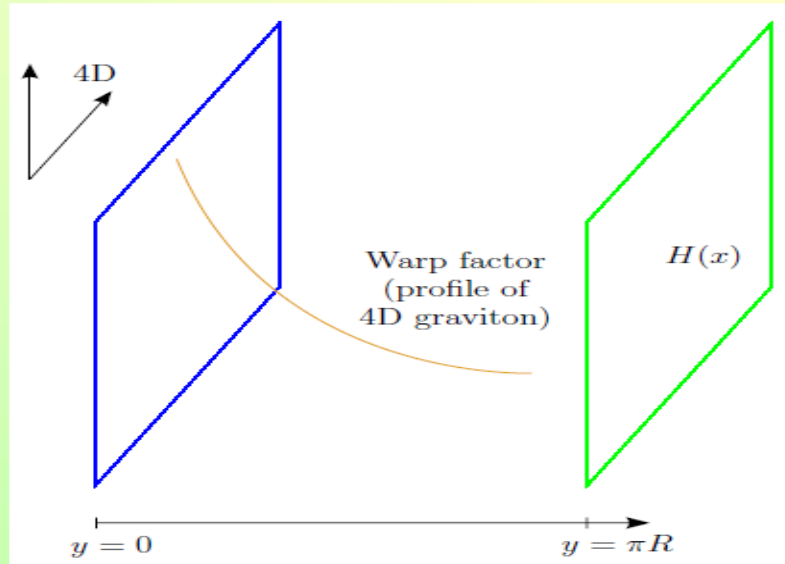
$$p^2 = p_5^2 = \frac{n^2}{R^2}, \quad (n = 0, 1, 2, \dots), \quad \Rightarrow \quad \text{KK Tower !}$$

➤ Masses Generation by **KK Compactification**:

—————▶ **Geometric “Higgs” Mechanism**

Warped 5d: Geometric “Higgs” Mechanism

➤ Warped 5d RS1 under S^1/Z_2 :



Randall-Sundrum

PRL.1999, hep-ph/9905221

➤ Warped 5d metric:

$$ds^2 = e^{2A(y)} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2,$$

$$ds^2 = e^{2A(z)} (\eta_{\mu\nu} dx^\mu dx^\nu + dz^2)$$

where

$$A(z) = -ky,$$

$$A(z) = -\ln(1 + kz),$$

$$k = \sqrt{-\Lambda/6}$$

$$\hat{g}_{MN} = e^{2A(z)} \hat{\eta}_{MN}$$

Geometric Higgs Mechanism for KK Gravitons

- Under S^1/Z_2 compactification, we impose the Boundary Conditions:

$$\partial_5 \hat{h}_{\mu\nu} \big|_{x^5=0,L} = 0, \quad \partial_5 \hat{\phi} \big|_{x^5=0,L} = 0, \quad \hat{h}_{\mu 5} \big|_{x^5=0,L} = 0$$

- 4d Quadratic KK Lagrangian contains Mixing Terms:

$$+ 2\mathbb{M}_n h_n \partial_\mu \mathcal{V}_n^\mu - 2\mathbb{M}_n h_n^{\mu\nu} \partial_\mu \mathcal{V}_{\nu,n} - \frac{3}{2} \mathbb{M}_n^2 h_n \phi_n - 3\mathbb{M}_n \partial_\mu \mathcal{V}_n^\mu \phi_n$$


which can be eliminated by **R_ξ gauge fixing**:

$$\begin{aligned} \mathcal{V}_n^\mu &= h_n^{\mu 5} \\ \phi_n &= h_n^{55} \end{aligned}$$

➔ $\mathcal{L}_{\text{GF}} = \int_0^L dz \hat{\mathcal{L}}_{\text{GF}} = - \sum_{n=0}^{\infty} \frac{1}{\xi_n} \left[(\mathcal{F}_n^\mu)^2 + (\mathcal{F}_n^5)^2 \right],$

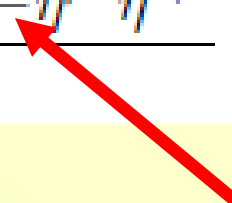
➔
$$\begin{aligned} \mathcal{F}_n^\mu &= \partial_\nu h_n^{\mu\nu} - \left(1 - \frac{1}{2\xi_n}\right) \partial^\mu h_n + \frac{1}{\sqrt{2}} \xi_n \mathbb{M}_n \mathcal{V}_n^\mu, \\ \mathcal{F}_n^5 &= \frac{1}{2} \mathbb{M}_n h_n - \sqrt{\frac{3}{2}} \xi_n \mathbb{M}_n \phi_n + \frac{1}{\sqrt{2}} \partial_\mu \mathcal{V}_n^\mu. \end{aligned}$$

➤ R_ξ Gauge Propagators:




$$\begin{aligned}
 \mathcal{D}_{nm}^{\mu\nu\alpha\beta}(p) = & -\frac{i\delta_{nm}}{2} \left\{ \frac{(\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta})}{p^2 + M_n^2} \right. \\
 & + \frac{1}{3} \left[\frac{1}{p^2 + M_n^2} - \frac{1}{p^2 + (3\xi_n - 2)M_n^2} \right] \left(\eta^{\mu\nu} - \frac{2p^\mu p^\nu}{M_n^2} \right) \left(\eta^{\alpha\beta} - \frac{2p^\alpha p^\beta}{M_n^2} \right) \\
 & + \frac{1}{M_n^2} \left[\frac{1}{p^2 + M_n^2} - \frac{1}{p^2 + \xi_n M_n^2} \right] (\eta^{\mu\alpha} p^\nu p^\beta + \eta^{\mu\beta} p^\nu p^\alpha + \eta^{\nu\alpha} p^\mu p^\beta + \eta^{\nu\beta} p^\mu p^\alpha) \\
 & \left. + \frac{4p^\mu p^\nu p^\alpha p^\beta}{\xi_n M_n^4} \left(\frac{1}{p^2 + \xi_n^2 M_n^2} - \frac{1}{p^2 + \xi_n M_n^2} \right) \right\}, \quad (2.4)
 \end{aligned}$$

- Feynman Gauge ($\xi_n=1$) Propagators take very simple form:


$$\mathcal{D}_{nm}^{\mu\nu\alpha\beta}(p) = -\frac{i\delta_{nm}}{2} \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2 + M_n^2}$$


- We also derive Unitary Gauge ($\xi_n=\infty$) Propagator:

$$\mathcal{D}_{nm,\text{UG}}^{\mu\nu\alpha\beta}(p) = -\frac{i\delta_{nm}}{2} \frac{\bar{\eta}^{\mu\alpha}\bar{\eta}^{\nu\beta} + \bar{\eta}^{\mu\beta}\bar{\eta}^{\nu\alpha} - \frac{2}{3}\bar{\eta}^{\mu\nu}\bar{\eta}^{\alpha\beta}}{p^2 + M_n^2}$$



$$\text{where } \bar{\eta}^{\mu\nu} = \eta^{\mu\nu} + p^\mu p^\nu / M_n^2.$$

➤ **KK Graviton propagator in $M_n \rightarrow 0$ limit:**

$$\begin{aligned} \mathcal{D}_{nm}^{\mu\nu\alpha\beta}(p) &= -\frac{i\delta_{nm}}{2} \left[\frac{(\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta})}{p^2} - \frac{1-\xi_n}{p^4} (\eta^{\mu\alpha}p^\nu p^\beta + \eta^{\mu\beta}p^\nu p^\alpha + \eta^{\nu\alpha}p^\mu p^\beta \right. \\ &\quad \left. + \eta^{\nu\beta}p^\mu p^\alpha - 2\eta^{\mu\nu}p^\alpha p^\beta - 2\eta^{\alpha\beta}p^\mu p^\nu) - 4(1-\xi_n)^3 \frac{p^\mu p^\nu p^\alpha p^\beta}{p^6} \right], \\ &= -\frac{i\delta_{nm}}{2} \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2}, \quad (\text{for } \xi_n=1). \end{aligned}$$


➔➔ It is **free from vDVZ (van-Dam-Veltman-Zakharov) Discontinuity !**


➤ **Compared with conventional massless graviton propagator:**

$$\begin{aligned} \mathcal{D}_{00}^{\mu\nu\alpha\beta}(p) &= -\frac{i}{2} \left[\frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2} + (\xi-1) \frac{\eta^{\mu\alpha}p^\nu p^\beta + \eta^{\mu\beta}p^\nu p^\alpha + \eta^{\nu\alpha}p^\mu p^\beta + \eta^{\nu\beta}p^\mu p^\alpha}{p^4} \right] \\ &= -\frac{i}{2} \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2}, \quad (\text{for } \xi=1). \end{aligned}$$



KK Graviton Propagator vs vDVZ Discontinuity

- **KK Graviton propagator in $M_n \rightarrow 0$ limit:**

$$\mathcal{D}_{nm}^{\mu\nu\alpha\beta}(p) = -\frac{i\delta_{nm}}{2} \left[\frac{(\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta})}{p^2} - \frac{1-\xi_n}{p^4} (\eta^{\mu\alpha}p^\nu p^\beta + \eta^{\mu\beta}p^\nu p^\alpha + \eta^{\nu\alpha}p^\mu p^\beta + \eta^{\nu\beta}p^\mu p^\alpha - 2\eta^{\mu\nu}p^\alpha p^\beta - 2\eta^{\alpha\beta}p^\mu p^\nu) - 4(1-\xi_n)^3 \frac{p^\mu p^\nu p^\alpha p^\beta}{p^6} \right],$$

$$= -\frac{i\delta_{nm}}{2} \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2}, \quad (\text{for } \xi_n=1).$$


- **Compared with Fierz-Pauli massive graviton propagator:**

$$\mathcal{D}_{\text{PF}}^{\mu\nu\alpha\beta}(p) = -\frac{i}{2} \frac{\bar{\eta}^{\mu\alpha}\bar{\eta}^{\nu\beta} + \bar{\eta}^{\mu\beta}\bar{\eta}^{\nu\alpha} - \frac{2}{3}\bar{\eta}^{\mu\nu}\bar{\eta}^{\alpha\beta}}{p^2 + M^2}$$


where $\bar{\eta}^{\mu\nu} = \eta^{\mu\nu} + p^\mu p^\nu / M^2$.

It has vDVZ discontinuity in $M_n \rightarrow 0$ limit!

KK Graviton Propagator vs vDVZ Discontinuity

- 5d KK Graviton **conserves** physical d.o.f in $M_n \rightarrow 0$ limit:

$$\lambda = \begin{matrix} 5 \\ 0, \pm 1, \pm 2 \end{matrix} = \begin{matrix} 2 \\ \pm 2 \end{matrix} + \begin{matrix} 2 \\ \pm 1 \end{matrix} + \begin{matrix} 1 \\ 0 \end{matrix}$$

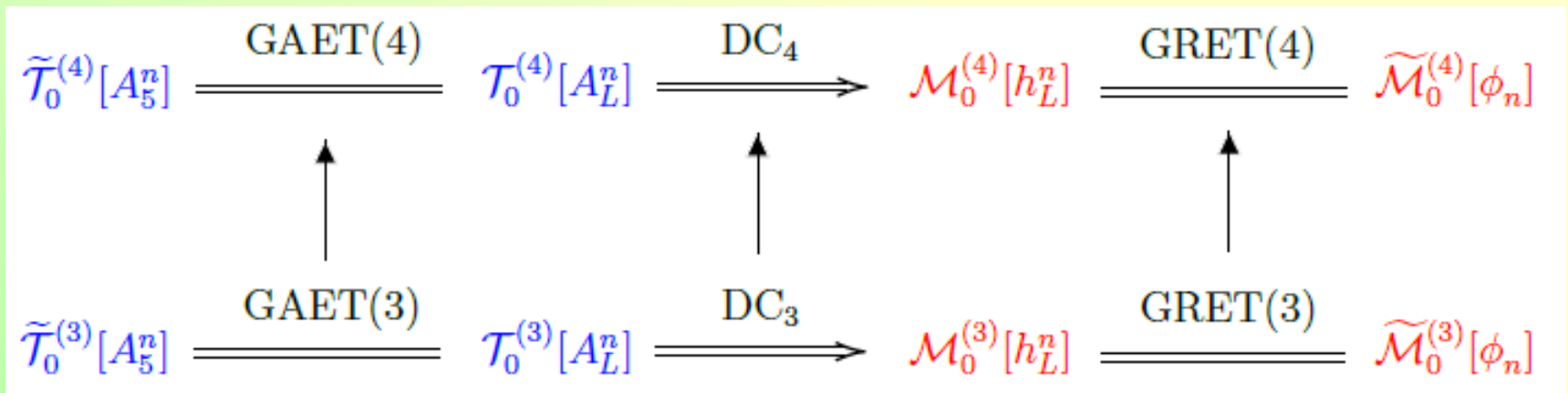
where $\lambda = 0, \pm 1$ d.o.f arise from 3 KK Goldstones $(h_n^{\mu 5}, h_n^{55})$.

- But, the **Massive Graviton** of Fierz-Pauli gravity does not conserve physical d.o.f in $M_n \rightarrow 0$ limit:

$$5 \neq 2$$

Schematic Summary:

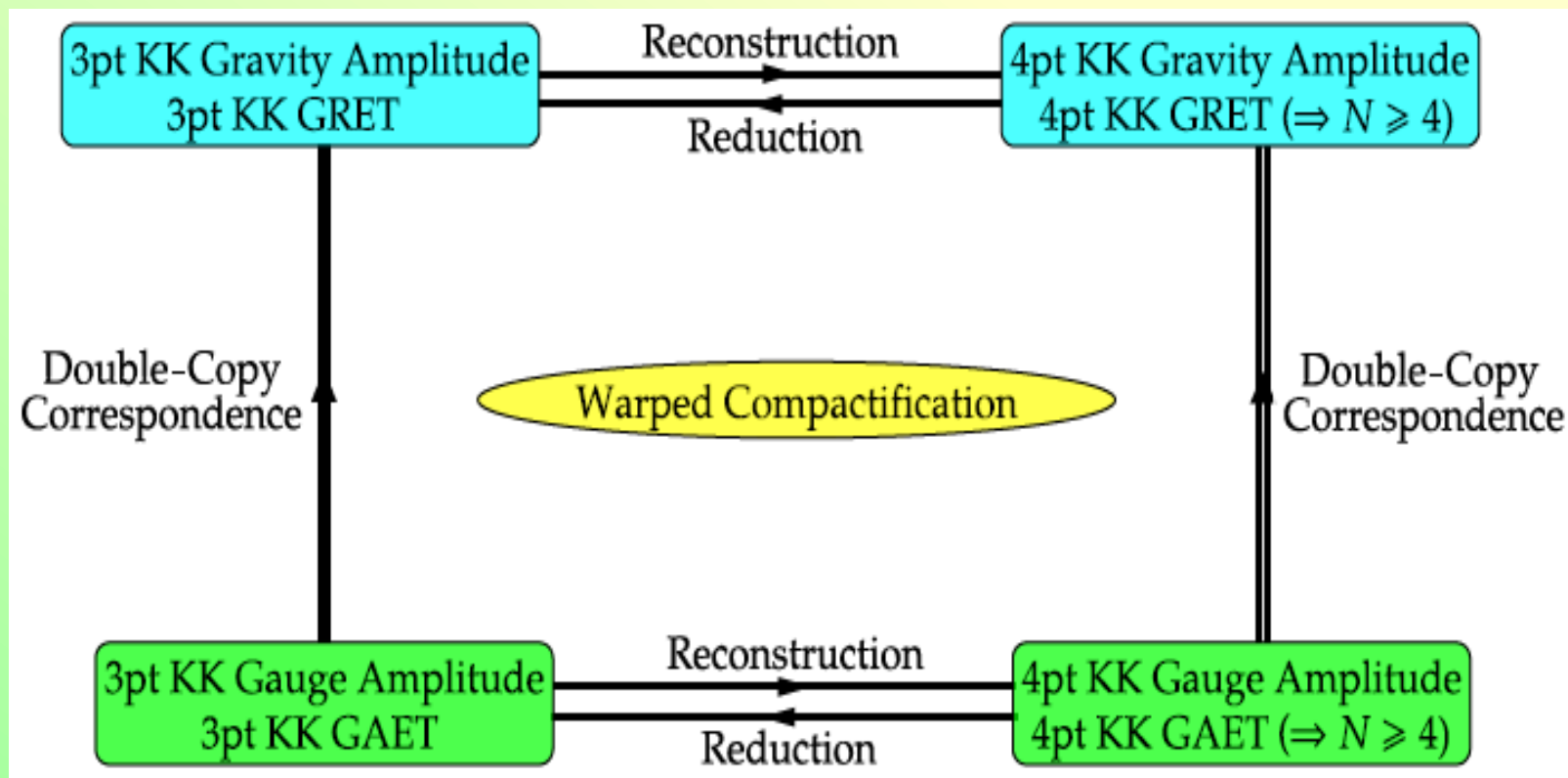
ET vs Double-Copy and from 3pt to 4pt Amplitudes



➤ This can be further extended to the case of $N > 4$ amplitudes.

KK ET & Double-Copy Correspondance: 3-Point to 4-Point

- Equivalence Theorem and Double-Copy Correspondence
from 3pt KK amplitudes to 4pt KK amplitudes and
from massive KK gauge amplitudes to massive KK gravitational amplitudes.



Geometric Higgs Mechanism via GRET

➤ Gravitational Equivalence Theorem (**GRET type-I**):

$$\mathcal{M}[h_{n_1}^L, \dots, h_{n_N}^L, \Phi] = C_{\text{mod}}^{n_j m_j} \mathcal{M}[\phi_{m_1}, \dots, \phi_{m_N}, \Phi] + O(\mathbb{M}_n/E_n\text{-suppressed})$$

$$C_{\text{mod}}^{n_j m_j} = C_{n_1 m_1} \cdots C_{n_N m_N} = \delta_{n_1 m_1} \cdots \delta_{n_N m_N} + O(\text{loop})$$

➤ Gravitational Equivalence Theorem (**GRET type-II**):

$$\mathcal{M}[h_{n_1}^{\pm 1}, \dots, h_{n_N}^{\pm 1}, \Phi] = \hat{C}_{\text{mod}}^{n_j m_j} \mathcal{M}[\mathcal{V}_{m_1}^{\pm 1}, \dots, \mathcal{V}_{m_N}^{\pm 1}, \Phi] + O(\mathbb{M}_n/E_n\text{-suppressed})$$

$$\hat{C}_{\text{mod}}^{n_j m_j} = \hat{C}_{n_1 m_1} \cdots \hat{C}_{n_N m_N} = (-i)^N \delta_{n_1 m_1} \cdots \delta_{n_N m_N} + O(\text{loop})$$

Generalized Power Counting Rule for KK Theories

➤ Generalized Energy Power Counting for KK Theories:

$$D_E = 2\mathcal{E}_{h_L} + (2L+2) + \sum_j \mathcal{V}_j (d_j - 2 + \frac{1}{2}f_j) .$$

➤ E-Power Counting of helicity-0 KK Graviton/Goldstone Amplitudes:

$$\begin{aligned} D_E(Nh_n^L) &= 2(N+L+1), & D_E(N\phi_n) &= 2(L+1), \\ \implies D_E(Nh_n^L) - D_E(N\phi_n) &= 2N, \end{aligned}$$

1

➤ E-Power Counting of helicity-1 KK Graviton/Goldstone Amplitudes:

$$\begin{aligned} D_E(Nh_n^{\pm 1}) &= N+2(L+1), & D_E(N\mathcal{V}_n^{\pm 1}) &= 2(L+1), \\ \implies D_E(Nh_n^{\pm 1}) - D_E(N\mathcal{V}_n^{\pm 1}) &= N. \end{aligned}$$

➤ E-Power Counting of KK Gauge/Goldstone boson Amplitudes:

$$\begin{aligned} D_E(NA_L^{an}) &= 4, & D_E(NA_5^{an}) &= 4 - N - \bar{V}_3^{\min}, \\ \implies D_E(NA_L^{an}) - D_E(NA_5^{an}) &= N + \bar{V}_3^{\min} \end{aligned}$$

➤ **Gravitational Equivalence Theorem (GET) identity:**

$$\mathcal{M}[h_{n_1}^L(k_1), \dots, h_{n_N}^L(k_N), \Phi] = \mathcal{M}[\phi_{n_1}(k_1), \dots, \phi_{n_N}(k_N), \Phi] + \sum_{1 \leq j \leq N} \mathcal{M}[\{\tilde{\Delta}_{n_j}, \phi_n\}, \Phi],$$

➤ **Energy Power Counting:**

$$D_E[Nh_n^L] = 2(N+1) + 2L,$$

$$D_E[N\phi_n] = 2 + 2L.$$

$$D_E[N\tilde{v}_n] = 2 + 2L,$$

➤ **We deduce h_L Amplitude has Large E-Cancellations:**



$$D_E[Nh_n^L] - D_E[N\phi_n] = 2N$$

➤ **For $N = 4$ scattering at tree level, GET proves Large E-cancellation:**



$$\mathbf{E}^{10} \rightarrow \mathbf{E}^2 \quad (\text{by } 10 - 2 = 8 \text{ powers})$$

KK Gauge Amplitudes & E-Cancellations at 3-Point

- 3pt Longitudinal KK gauge boson amplitude ($A_L A_L A_T$) and KK Goldstone amplitude ($A_5 A_5 A_T$):

$$\mathcal{T}_0[A_L^{an_1} A_L^{bn_2} A_{\pm}^{cn_3}] = -ig f^{abc} (\epsilon_3 \cdot p_1) \frac{(M_{n_3}^2 - M_{n_1}^2 - M_{n_2}^2)}{M_{n_1} M_{n_2}} a_{n_1 n_2 n_3},$$
$$\tilde{\mathcal{T}}_0[A_5^{an_1} A_5^{bn_2} A_{\pm}^{cn_3}] = -i2g f^{abc} (\epsilon_3 \cdot p_1) \tilde{a}_{n_1 n_2 n_3}$$

- The most fundamental GAET:

→
$$\mathcal{T}_0[A_L^{an_1} A_L^{bn_2} A_{\pm}^{cn_3}] = -\tilde{\mathcal{T}}_0[A_5^{an_1} A_5^{bn_2} A_{\pm}^{cn_3}]$$

enforcing energy cancellation: $E^3 \rightarrow E^1$

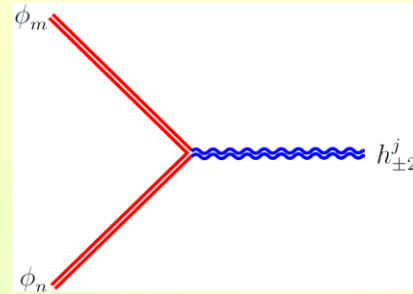
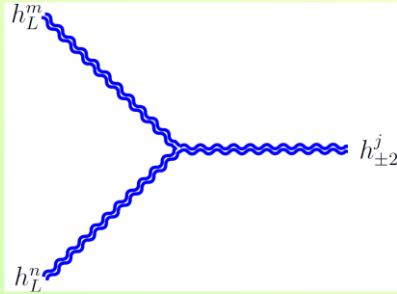
and it requires the coupling-mass condition:

→
$$(M_{n_1}^2 + M_{n_2}^2 - M_{n_3}^2) a_{n_1 n_2 n_3} = 2M_{n_1} M_{n_2} \tilde{a}_{n_1 n_2 n_3}$$

from which all N-point KK Amplitudes and GAET (N>3) can be derived !

KK Graviton Amplitudes & E-Cancellations at 3-Point

- 3pt Longitudinal KK graviton amplitude ($h_L h_L h_T$) and KK Goldstone amplitude ($\phi_n \phi_m h_T$) :



$$\mathcal{M}_0[h_{n_1}^L h_{n_2}^L h_{n_3}^{\pm 2}] = \frac{\kappa(\epsilon_3 \cdot p_1)^2}{6M_{n_1}^2 M_{n_2}^2} \left[2M_{n_1}^2 M_{n_2}^2 + (M_{n_1}^2 + M_{n_2}^2 - M_{n_3}^2)^2 \right] \alpha_{n_1 n_2 n_3},$$

$$\tilde{\mathcal{M}}_0[\phi_{n_1} \phi_{n_2} h_{n_3}^{\pm 2}] = \kappa(\epsilon_3 \cdot p_1)^2 \tilde{\beta}_{n_1 n_2 n_3},$$

- E Cancellations for above 3-point KK Graviton Amplitude: $E^6 \rightarrow E^2$

Gauge-Gravity Duality: Double-Copy

Massless vs Massive

Li, Hang, HJH, JHEP (2022), 2209.11191
Hang, Zhao, HJH, Qiu, JHEP (2025), 2406.12713
Hang, HJH, PRD (2022), 2106.04568, 2207.11214
Li, Hang, HJH, He, JHEP(2022), 2111.12042
Hang, HJH, Shen, JHEP(2022), 2110.05399
Research (2023), 2406.13671
H.X. Liu, Z.X. Yi, HJH, **2512.10870 [hep-th]** (66pp)

$$(\text{GR}) = (\text{QCD})^2$$

$$(\text{Gravity}) = (\text{Gauge Theory})^2$$

$$(\text{引力}) = (\text{规范力})^2$$

Double-Copy: KLT vs BCJ

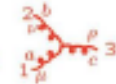
- for Amplitudes of **Massless open/closed strings** and
for **massless gauge/graviton bosons**:

1985: Kawai, Lewellen, Tye (KLT): “closed string amp=open-string amp²”

Field-theory limit:



Yang-Mills
gauge theory:



$$-gf^{abc}(\eta_{\mu\nu}(k_1 - k_2)_\rho + \text{cyclic})$$

Einstein
gravity:



$$i\kappa(\eta_{\mu\nu}(k_1 - k_2)_\rho + \text{cyclic}) \\ \times (\eta_{\alpha\beta}(k_1 - k_2)_\gamma + \text{cyclic})$$

“square” of
Yang-Mills
vertex.

↙ gravity

↙ gauge theory color ordered

$$M_4^{\text{tree}}(1,2,3,4) = -i s_{12} A_4^{\text{tree}}(1,2,3,4) A_4^{\text{tree}}(1,2,4,3)$$

2008: Bern, Carrasco, Johansson (BCJ):

$$\mathcal{A}_4^{\text{tree}} = g^2 \left(\frac{n_s c_s}{s} + \frac{n_t c_t}{t} + \frac{n_u c_u}{u} \right)$$

$$\mathcal{A}_4^{\text{tree}} \Big|_{c_i \rightarrow n_i} \equiv -i \mathcal{M}_4^{\text{tree}} = \frac{n_s^2}{s} + \frac{n_t^2}{t} + \frac{n_u^2}{u}$$



$$c_s + c_t + c_u = 0$$

$$n_s + n_t + n_u = 0$$

Massive Double-Copy vs Mass Spectral Condition

- For the massive KK graviton amplitude:

$$\mathcal{M}\left[h_{\sigma_1}^{n_1} h_{\sigma_2}^{n_2} h_{\sigma_3}^{n_3} h_{\sigma_4}^{n_4}\right] = -\frac{\kappa^2}{16} \sum_j \sum_{\lambda_k, \lambda'_k} \left(\prod_{k=1}^4 C_{\lambda_k \lambda'_k}^{\sigma_k} \right) \frac{\mathcal{N}_j^{\text{P}}(\lambda_k) \mathcal{N}_j^{\text{P}}(\lambda'_k)}{s_j - M_{\text{nn}_j}^2},$$

we impose the Generalized Gauge Transformation:

$$\mathcal{N}_j^{\text{P}'} = \mathcal{N}_j^{\text{P}} + (s_j - M_{\text{nn}_j}^2) \times \Delta,$$

from which we deduce conditions:

$$\sum_j \mathcal{N}_j^{\text{P}} = 0, \quad \sum_j (s_j - M_{\text{nn}_j}^2) = 0.$$

- This leads to the 4-point **Mass Spectral Condition**: ➔ **Nontrivial !**
 ➔ **Does Not always hold !**

$$\sum_{i=1}^4 M_{\text{n}_i}^2 = M_{\text{nn}_s}^2 + M_{\text{nn}_t}^2 + M_{\text{nn}_u}^2$$

Massive Double-Copy vs Mass Spectral Condition

- Start from a general 4-point **Mass Spectral Condition**:

$$M_1^2 + M_2^2 + M_3^2 + M_4^2 = M_{12}^2 + M_{13}^2 + M_{14}^2$$

- We ask: **How to solve it and What is the solution ??**

- For the **flat 5d Toroidal Compactification** of S^1 , we have:

$$n_1^2 + n_2^2 + n_3^2 + n_4^2 = (n_1 + n_2)^2 + (n_1 + n_3)^2 + (n_1 + n_4)^2$$

which leads to the condition:

$$n_1 + n_2 + n_3 + n_4 = 0$$

- This is just the **KK number (5d momentum) Conservation!**
- E.g., it does not hold for **5d Orbifold Compactification** or **5d Warped Compactification** ! (additional treatments needed.)

Massive Double-Copy vs Mass Spectral Condition

- Start from a general 4-point Mass Spectral Condition:

$$M_1^2 + M_2^2 + M_3^2 + M_4^2 = M_{12}^2 + M_{13}^2 + M_{14}^2$$

- We ask: How to solve it and What is the solution ???

- In fact, starting with 3 rather modest conditions:

- 1). A massive theory contains at least 2 types of particles with Unequal masses.
- 2). There exists only a simple pole in each of (s, t, u) channels.
- 3). Each scattering amplitude should include the contributions from all 3 kinematic channels of (s, t, u).

we can prove:

→→ **Toroidal Compactification of flat Extra dimensions** is the **Unique** Solution !

Massive Double-Copy vs Mass Spectral Condition

- General 4-point **Mass Spectral Condition**:

$$M_1^2 + M_2^2 + M_3^2 + M_4^2 = M_{12}^2 + M_{13}^2 + M_{14}^2 \quad (1)$$

- **Proof of the Unique Solution**:

We identified the group structure underlying condition (1) is a **product of Integer Groups \mathbb{Z}^r** (with rank r) in the **Finitely Generated Abelian Group**:

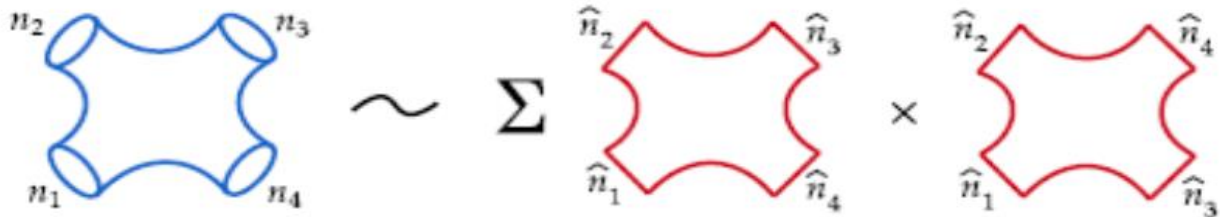
$$\mathcal{G} \cong \mathbb{Z}^r \oplus \mathbb{Z}_{p_1} \oplus \mathbb{Z}_{p_2} \oplus \cdots \oplus \mathbb{Z}_{p_s}, \quad (r, s \in \mathbb{N})$$

- Then, we proved that the unique solution to Condition (1) is given by $M_{\{k\}}^2 = k^2 M_{\{1\}}^2$, for the case $r = 1$ and by $M_n^2 = n \bar{M}^2 n^T$ for the case $r > 1$.

- By inspecting all the known consistent QFTs, we conclude that **only the KK theories with δ ($= r$) extra dimensions under the toroidal compactification** could have a mass spectrum behave exactly as above. These theories also hold the **KK number conservation**.

Gauge Forces vs Gravity Force

- Massive Case: $GR = (\text{Gauge Theory})^2 \quad ??$
- What happens to **massive Kaluza-Klein Theory** ??
- **First Principle Approach:** Using KK bosonic string theory, we derived Massive KLT Relations between product of KK Open String Amplitudes and KK Closed String Amplitude. Taking Field Theory Limit, we derived Massive KLT Relations between product of KK Gauge Boson Amplitudes and KK Graviton Amplitude:



$$\mathcal{M}[1_L^n 2_L^n 3_L^n 4_L^n] = \frac{\kappa^2}{32} \sum_{\{a_j, b_j\}} \hat{\rho}_{ab} \left\{ (s - 4M_n^2) \mathcal{T}_{a_j}[1^{+n} 2^{+n} 3^{-n} 4^{-n}] \mathcal{T}_{b_j}[1^{+n} 2^{+n} 4^{-n} 3^{-n}] \right. \\ \left. + s \mathcal{T}_{a_j}[1^{+n} 2^{-n} 3^{+n} 4^{-n}] \mathcal{T}_{b_j}[1^{+n} 2^{-n} 4^{+n} 3^{-n}] \right. \\ \left. + s \mathcal{T}_{a_j}[1^{+n} 2^{-n} 3^{-n} 4^{+n}] \mathcal{T}_{b_j}[1^{+n} 2^{-n} 4^{-n} 3^{+n}] \right\}$$

Topological Mass Generation & Scattering Amplitudes Topological Equivalence Theorem & Double-Copy for Chern-Simons Gauge & Gravity Theories

H.X. Liu, Z.X. Yi, HJH, [arXiv:2512.10870 \[hep-th\]](#) (66 pp)
Hang, HJH, Shen JHEP(2022), 2110.05399
Research (2023), 2406.13671

3d Topological Massive Gauge & Gravity Theories

- 3d Spacetime has distinctive features characterized by **Gauge & Gravitational Chern-Simons terms**, which generate **topological masses** for gauge bosons and gravitons, realize **fractional statistics** and predict the existence of **anyon-like** quasi-particles.
- Chern-Simons (CS) terms are topological invariants in mathematics and serve as the theoretical key ingredients of a wide range of applications in modern physics, including fractional quantum Hall effect, models of high-temperature superconductivity and strongly correlated systems, & topological quantum computing.
- Studying structure of scattering amplitudes of massive gauge bosons and gravitons in the Chern-Simons theories provides an important means for understanding the mechanism of **Topological Mass Generations** and for realizing Gauge-Gravity duality connection: **$(\text{Gravity}) = (\text{Gauge})^2$**

3d Topological Massive Gauge & Gravity Theories

➤ 3d Chern-Simons (CS) Topological Massive Gauge Theories:

$$\mathcal{L}_{\text{TMQED}} = -\frac{1}{4}F_{\mu\nu}^2 + \frac{1}{2}\tilde{m}\varepsilon^{\mu\nu\rho}A_\mu\partial_\nu A_\rho,$$

$$\mathcal{L}_{\text{TMYM}} = -\frac{1}{2}\text{tr}F_{\mu\nu}^2 + \tilde{m}\varepsilon^{\mu\nu\rho}\text{tr}\left(A_\mu\partial_\nu A_\rho - \frac{i2g}{3}A_\mu A_\nu A_\rho\right),$$

➤ CS mass is geometrized and is related to CS Level:

$$n = 4\pi\tilde{m}/g^2 \in \mathbb{Z}$$

➤ 3d CS Topological Massive Gravity Theories:

$$S_{\text{TMG}} = -\frac{2}{\kappa^2}\int d^3x\left[\sqrt{-g}R - \frac{1}{2\tilde{m}}\varepsilon^{\mu\nu\rho}\Gamma^\alpha_{\rho\beta}\left(\partial_\mu\Gamma^\beta_{\alpha\nu} + \frac{2}{3}\Gamma^\beta_{\mu\gamma}\Gamma^\gamma_{\nu\alpha}\right)\right]$$

E³

E²

Topological Mass Generation

- Conversion of physical d.o.f between $m = 0$ and $m \neq 0$:

$$A_P^a = \frac{1}{\sqrt{2}} (A_T^a + A_L^a) \longrightarrow A_T^a$$

- 2 orthogonal unphysical states:

$$\begin{aligned} A_X^a &= \epsilon_X^\mu A_\mu^a = \frac{1}{\sqrt{2}} (A_T^a - A_L^a), \\ A_S^a &= \epsilon_S^\mu A_\mu^a, \end{aligned}$$

- The physical d.o.f is conserved under $m \rightarrow 0$ limit:

$$1 = 1$$

- Then, we ask:

$$\mathcal{T}[A_P^{a_1}, \dots, A_P^{a_N}, \Phi] \stackrel{?}{=} \mathcal{T}[\tilde{A}_T^{a_1}, \dots, \tilde{A}_T^{a_N}, \Phi]$$

Topological Mass Generation: TET

- Indeed we can derive a **new identity** to connect the Scattering Amplitudes:

$$\mathcal{T}[A_{\text{P}}^{a_1}, \dots, A_{\text{P}}^{a_N}, \Phi] = \mathcal{T}[\tilde{A}_{\text{T}}^{a_1}, \dots, \tilde{A}_{\text{T}}^{a_N}, \Phi] + \mathcal{T}_v,$$
$$\mathcal{T}_v = \sum_{j=1}^N \mathcal{T}[\tilde{v}^{a_1}, \dots, \tilde{v}^{a_j}, \tilde{A}_{\text{T}}^{a_{j+1}}, \dots, \tilde{A}_{\text{T}}^{a_N}, \Phi],$$

- Under high energy expansion, we derive, **at S-matrix level**,
Topological Equivalence Theorem (TET):

$$\mathcal{T}[A_{\text{P}}^{a_1}, \dots, A_{\text{P}}^{a_N}, \Phi] = \mathcal{T}[\tilde{A}_{\text{T}}^{a_1}, \dots, \tilde{A}_{\text{T}}^{a_N}, \Phi] + \mathcal{O}\left(\frac{m}{E}\right),$$

➤ **E Cancellations: $E^4 \rightarrow E^1$**

Amplitude	$\times \bar{s}_0^2$	$\times \bar{s}_0^{3/2}$	$\times \bar{s}_0$
\mathcal{M}_s	$-\frac{99+28c_{2\theta}+c_{4\theta}}{1-c_{2\theta}}$	$-i14(15c_\theta+c_{3\theta})\csc\theta$	$-\frac{2(75+326c_{2\theta}+47c_{4\theta})}{1-c_{2\theta}}$
\mathcal{M}_t	$\frac{99+28c_{2\theta}+c_{4\theta}}{4(1-c_\theta)}$	$i(102+105c_\theta+70c_{2\theta}+7c_{3\theta}+4c_{4\theta})\csc\theta$	$\frac{75-107c_\theta+326c_{2\theta}+268c_{3\theta}+47c_{4\theta}+31c_{5\theta}}{1-c_{2\theta}}$
\mathcal{M}_u	$\frac{99+28c_{2\theta}+c_{4\theta}}{4(1+c_\theta)}$	$i(-102+105c_\theta-70c_{2\theta}+7c_{3\theta}-4c_{4\theta})\csc\theta$	$\frac{75+107c_\theta+326c_{2\theta}-268c_{3\theta}+47c_{4\theta}-31c_{5\theta}}{1-c_{2\theta}}$
Sum	0	0	0

➤ **Under high energy expansion:**

$$\mathcal{M}_0[4h_p] = -\frac{i\kappa^2}{2048} m s_0^{\frac{1}{2}} (494c_\theta + 19c_{3\theta} - c_{5\theta}) \csc^3\theta$$

➤ **Energy Cancellations for 4 graviton scattering amplitudes:**

$$E^4 \rightarrow E^0 \quad \longrightarrow \quad \mathcal{O}(E^{12}) \longrightarrow \mathcal{O}(E^1), \quad (\text{for } \mathcal{E}_{h_p}=4 \text{ in 3d TMG})$$

Topological Mass Generation: Covariant TMG

➤ Introduce dilaton field via conformal factor:

$$g_{\mu\nu} = \bar{g}_{\mu\nu} e^{-\kappa\phi},$$


$$\mathcal{L}_{\text{TMG}}^\phi = \frac{-1}{\kappa^2} \left\{ \sqrt{-\bar{g}} e^{-\kappa\phi/2} \left[R + \frac{\kappa^2}{2} \bar{g}_{\mu\nu} \partial^\mu \phi \partial^\nu \phi \right] - \frac{\varepsilon^{\mu\nu\rho}}{2\tilde{m}} \Gamma^\alpha_{\rho\beta} \left(\partial_\mu \Gamma^\beta_{\alpha\nu} + \frac{2}{3} \Gamma^\beta_{\mu\gamma} \Gamma^\gamma_{\nu\alpha} \right) \right\}$$

➤ Construct Gauge-fixing Terms:

$$\begin{aligned} \mathcal{L}_{\text{TMG}}^{\text{GF1}} &= \frac{1}{2\xi} (\mathcal{F}_{\text{GF1}}^\mu)^2, & \mathcal{F}_{\text{GF1}}^\mu &= \partial_\nu \bar{h}^{\mu\nu} - \frac{1}{2} \partial^\mu (\bar{h} - \xi\phi), \\ \mathcal{L}_{\text{TMG}}^{\text{GF2}} &= \frac{1}{2\zeta} (\mathcal{F}_{\text{GF2}}^\mu)^2, & \mathcal{F}_{\text{GF2}}^\mu &= \frac{1}{2} \partial^\mu (\bar{h} - \zeta\phi), \end{aligned}$$

➤ Landau Gauge Propagators:

$$\begin{aligned} \mathcal{D}_{\mu\nu\alpha\beta}^h(p) &= -\frac{mp^\rho}{2p^2(p^2+m^2)} \left(\varepsilon_{\rho\mu\alpha} \bar{\eta}_{\nu\beta} + \varepsilon_{\rho\mu\beta} \bar{\eta}_{\nu\alpha} + \varepsilon_{\rho\nu\alpha} \bar{\eta}_{\mu\beta} + \varepsilon_{\rho\nu\beta} \bar{\eta}_{\mu\alpha} \right) \\ &\quad + \frac{im^2}{p^2(p^2+m^2)} \left(\bar{\eta}_{\mu\alpha} \bar{\eta}_{\nu\beta} + \bar{\eta}_{\mu\beta} \bar{\eta}_{\nu\alpha} - \bar{\eta}_{\mu\nu} \bar{\eta}_{\alpha\beta} \right), \\ D^\phi(p) &= -\frac{i}{p^2}, \end{aligned}$$

 1/p³

➤ Landau Gauge Propagators ($\zeta = 0$ and $\xi = 0$):

$$\begin{aligned} \mathcal{D}_{\mu\nu\alpha\beta}^h(p) &= -\frac{mp^\rho}{2p^2(p^2+m^2)} \left(\varepsilon_{\rho\mu\alpha}\bar{\eta}_{\nu\beta} + \varepsilon_{\rho\mu\beta}\bar{\eta}_{\nu\alpha} + \varepsilon_{\rho\nu\alpha}\bar{\eta}_{\mu\beta} + \varepsilon_{\rho\nu\beta}\bar{\eta}_{\mu\alpha} \right) \\ &\quad + \frac{im^2}{p^2(p^2+m^2)} \left(\bar{\eta}_{\mu\alpha}\bar{\eta}_{\nu\beta} + \bar{\eta}_{\mu\beta}\bar{\eta}_{\nu\alpha} - \bar{\eta}_{\mu\nu}\bar{\eta}_{\alpha\beta} \right), \\ D^\phi(p) &= -\frac{i}{p^2}, \end{aligned}$$

1/p³

➤ Unitary Gauge Propagator ($\zeta = \infty$ and $\xi = 0$):

$$\begin{aligned} D_{\mu\nu\alpha\beta}^{h(U)}(p) &= \frac{-i(\bar{\eta}_{\mu\alpha}\bar{\eta}_{\nu\beta} + \bar{\eta}_{\mu\beta}\bar{\eta}_{\nu\alpha} - \bar{\eta}_{\mu\nu}\bar{\eta}_{\alpha\beta})}{p^2+m^2} + \frac{i\xi}{p^4} (\eta_{\mu\alpha}p_\nu p_\beta + \eta_{\mu\beta}p_\nu p_\alpha + \eta_{\nu\alpha}p_\mu p_\beta + \eta_{\nu\beta}p_\mu p_\alpha) \\ &\quad - \frac{mp^\rho}{2p^2(p^2+m^2)} \left(\varepsilon_{\rho\mu\alpha}\bar{\eta}_{\nu\beta} + \varepsilon_{\rho\mu\beta}\bar{\eta}_{\nu\alpha} + \varepsilon_{\rho\nu\alpha}\bar{\eta}_{\mu\beta} + \varepsilon_{\rho\nu\beta}\bar{\eta}_{\mu\alpha} \right) \\ &\quad + \frac{i}{p^2} \left(\bar{\eta}_{\mu\alpha}\bar{\eta}_{\nu\beta} + \bar{\eta}_{\mu\beta}\bar{\eta}_{\nu\alpha} - 2\bar{\eta}_{\mu\nu}\bar{\eta}_{\alpha\beta} \right) - \frac{i2}{p^4} (\eta_{\mu\nu}p_\alpha p_\beta + \eta_{\alpha\beta}p_\mu p_\nu), \end{aligned}$$

1/p²

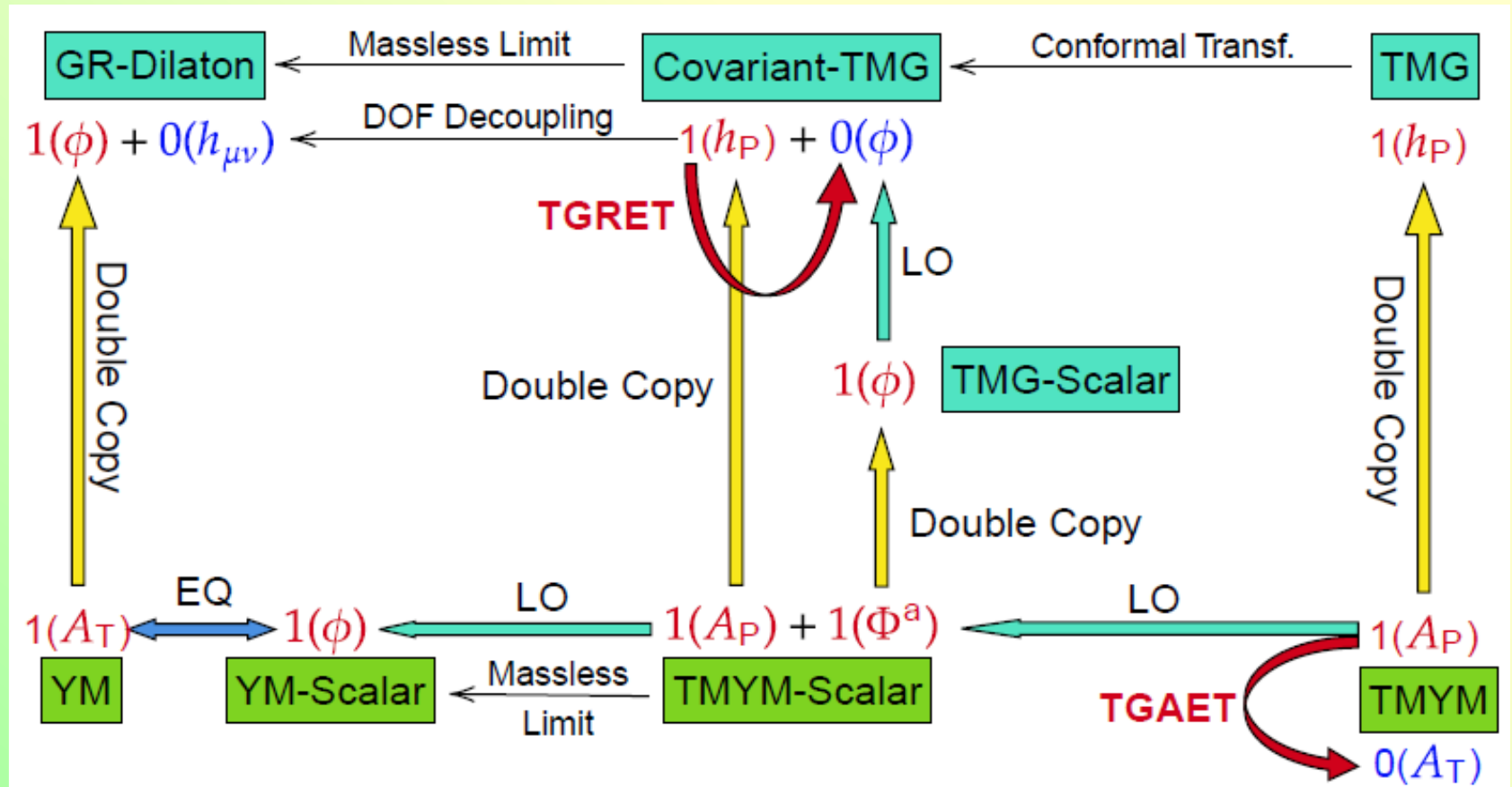
➤ **Topological Gravity Equivalence Theorem (TGRET) :**

$$\mathcal{M}[h_{\text{P}}(p_1), \dots, h_{\text{P}}(p_N), \Phi] = \left(\frac{1}{2}\right)^N C_{\text{mod}} \mathcal{M}[\phi(p_1), \dots, \phi(p_N), \Phi] + O\left(\frac{m}{E}\right),$$

➤ **In comparison with TGAET:**

$$\mathcal{T}[A_{\text{P}}^{a_1}(p_1), \dots, A_{\text{P}}^{a_N}(p_N), \Phi] = \left(\frac{1}{\sqrt{2}}\right)^N \mathcal{T}[A_{\text{T}}^{a_1}, \dots, A_{\text{T}}^{a_N}, \Phi] + O\left(\frac{m}{E}\right)$$

Topological Mass Generation: Summary



Generalized Power Counting for 3d Chern-Simons

- **Generalized Energy Power Counting** for 3d CS physical gauge boson scattering amplitudes:

$$D_E = (\mathcal{E}_{A_P} - \mathcal{E}_v) + (4 - \mathcal{E} - \overline{\mathcal{V}}_3) - L,$$

- **Large E Cancellations in N physical A_P amplitudes by TET:**

$$\Delta D_E = D_E[NA_P^a] - D_E[NA_T^a] = N$$

- **E-Power Counting of N physical Graviton Amplitudes:**

$$D_E[Nh_P] = 2N + 2 + (\mathcal{V} - I_h + L) = 2N + 3$$

$$D_E^U[Nh_P] = (2N + 2) + \mathcal{V} + L = (2N + 3) + I_h$$

$$D_{E(0)}[N\phi] = 3 - n = 3 - \frac{1}{2}N, \quad (N = 2n),$$

$$D_{E(0)}[N\phi] = 2 - n = \frac{1}{2}(5 - N), \quad (N = 2n + 1),$$

N=4 Scattering Amplitudes:

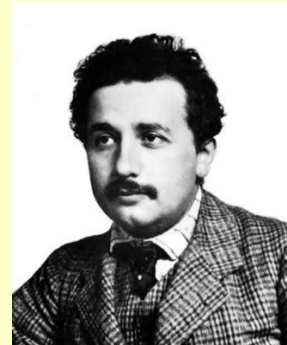
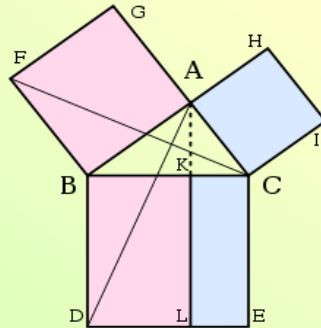
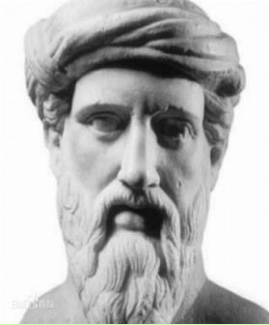
$E^{12} \rightarrow E^1$ (Unitary Gauge)

$E^{11} \rightarrow E^1$ (Landau Gauge)

$$D_{E(0)}[Nh_P] - D_{E(0)}[N\phi] = \begin{cases} \frac{5}{2}N, & (\text{for } N = \text{even}), \\ \frac{1}{2}(5N + 1), & (\text{for } N = \text{odd}), \end{cases}$$

$$D_{E(0)}^U[Nh_P] - D_{E(0)}[N\phi] = \begin{cases} \frac{7}{2}N - 3, & (\text{for } N = \text{even}), \\ \frac{1}{2}(7N - 5), & (\text{for } N = \text{odd}). \end{cases}$$

Deep Relations: Importance of Equality & Square Law



- [EG] Pythagoras Theorem: $a^2 + b^2 = c^2$ → Fermat Last Theorem: $a^n + b^n \neq c^n$ ($n > 2$)
- [SR] Mass-Energy Equivalence: $E = Mc^2$ and $P_\mu^2 = M^2 c^2$
- [GR] Inertial Mass vs Gravitational-Mass Equivalence: $M_G = M_I \rightarrow g = a$
- Gravity Force from Gauge Force: (valid for $M = 0$ and $M \neq 0$)

→
 $M = 0$
(Conventional works)

$$(\text{Gravity}) = (\text{Gauge Force})^2$$

←
 $M \neq 0$
(This Talk)



→ **Much more can be explored !**

Thank You