

# **Origin of Mass and Scattering Amplitudes: from Higgs to Pauli, Kaluza-Klein and Chern-Simons**

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Wuhan, December 15, 2025

# References

## ➤ Based on series of works:

You, Song, HJH, Han, PRD(2025), 2412.16033

Han, HJH, You, Song, PRD-Lett (2025), 2412.21045

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Y. Hang, W.W. Zhao, HJH, Y. Qiu, JHEP (2025), 2406.12713 [hep-th] (103 pp).

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Y. Li, Y.F. Hang, HJH, S. He, JHEP (2022), 2111.12042 [hep-th].

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Y.F. Hang, HJH, C. Shen, JHEP (2022), 2110.05399 [hep-th].

Y.F. Hang, HJH, C. Shen, Research (2023), 2406.13671 [hep-th].

## ➤ Early works:

Chivukula, Dicus, HJH, PLB (2002) [hep-ph/0111016]

Chivukula, HJH, PLB (2002) [hep-ph/0201164]

H.J. He, IJMP (2005), APS-2004 [hep-ph/0412113]

Chivukula, Dicus, HJH, Nandi, PLB (2003) [hep-ph /0302263]

## Making of the Standard Model (123)

$$\mathcal{L} = -\frac{1}{4g'^4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4g^2} W_{\mu\nu}^a W^{\mu\nu a} - \frac{1}{4g_s^2} G_{\mu\nu}^a G^{\mu\nu a} + \bar{Q}_i i \not{D} Q_i + \bar{u}_i i \not{D} u_i + \bar{d}_i i \not{D} d_i + \bar{L}_i i \not{D} L_i + \bar{\ell}_i i \not{D} \ell_i + \left( Y_u^{ij} \bar{Q}_i u_j \tilde{H} + Y_d^{ij} \bar{Q}_i d_j H + Y_l^{ij} \bar{L}_i \ell_j H + c.c. \right) - \lambda (H^\dagger H)^2 + \lambda v^2 H^\dagger H - (D^\mu H)^\dagger D_\mu H$$

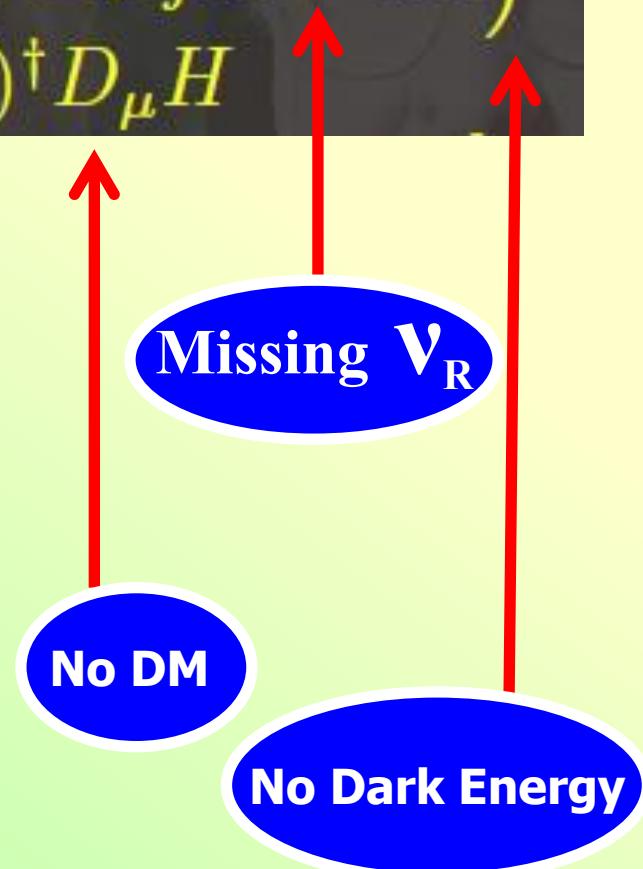
Is **SM** Structure complete? → **not really !!**

→ What is really **Missing** in the SM ???

1)  **$v_R$**  Within SM Structure **not yet found** !

2) **No Dark Matter** ( ? X )

3) **No Dark Energy** ( ? ✓ )



# Making of the Standard Model (123)

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4g'^4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{4g^2}W_{\mu\nu}^aW^{\mu\nu a} - \frac{1}{4g_s^2}G_{\mu\nu}^aG^{\mu\nu a} \\ & + \bar{Q}_i i\cancel{D}Q_i + \bar{u}_i i\cancel{D}u_i + \bar{d}_i i\cancel{D}d_i + \bar{L}_i i\cancel{D}L_i + \bar{\ell}_i i\cancel{D}\ell_i \\ & + \left( Y_u^{ij} \bar{Q}_i u_j \tilde{H} + Y_d^{ij} \bar{Q}_i d_j H + Y_l^{ij} \bar{L}_i \ell_j H + c.c. \right) \\ & - \lambda (H^\dagger H)^2 + \lambda v^2 H^\dagger H + (D^\mu H)^\dagger D_\mu H\end{aligned}$$

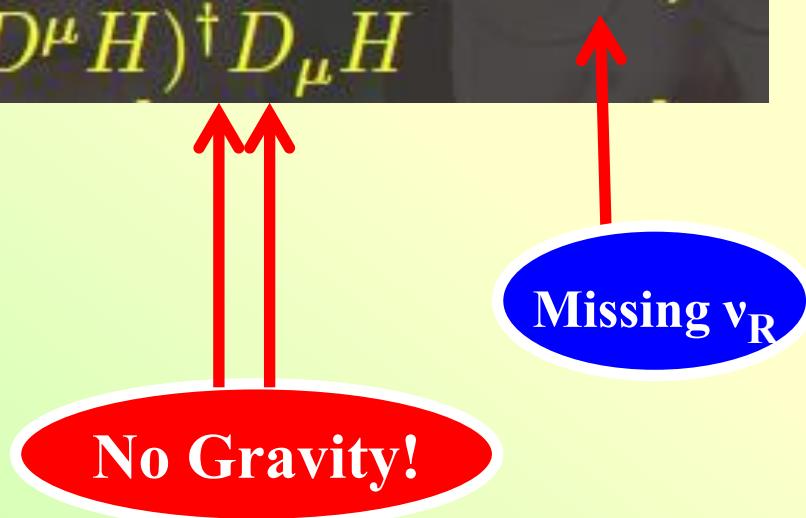
**SM Structure is NOT yet complete ....**

→ What is *Missing* in the SM ??

→→ No full understanding on Quantum Gravity

at both the Largest and Smallest Scales !!

→→ UV-IR correspondence ? ...



# **Origin of Neutrino Mass Generation and How to Probe it ?**

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**You, Song, HJH, Han, PRD(2025), 2412.16033  
Han, HJH, You, Song, PRD-Lett (2025), 2412.21045**

# Wolfgang Pauli: Father of the Neutrinos



W. Pauli (1900-1958)

- 1914: James Chadwick discovered that  $\beta$  decay has continuous electron energy spectrum and could not explain it with 2-body final state.
- In 1929, Niels Bohr attempted to explain this by giving up energy conservation (X).
- 1930: Pauli postulated that the final state is 3-body, containing a new massless particle with electric charge 0 & spin-1/2. This naturally explains the continuous E-spectrum + conserves energy. ( ✓ )
- 1932: Fermi named this particle as “Neutrino”.
- 1933: Fermi wrote down 4-Fermion Interaction to describe  $\beta$  decays in weak interaction.



E. Fermi (1901-1954)

## Last Missing Piece of the SM

- For conventional SM setup before 1998, neutrinos were assumed for simplicity to be massless and have only left-handed components because the SM is structured to have all the right-handed fermions be weak singlets in each fermion family, where the right-handed neutrinos ( $v_R$ ) are pure gauge singlets and their absence does not affect the gauge anomaly cancellation of the SM.
- Weinberg realized that without  $v_R$ , the left-handed neutrinos can acquire small Majorana masses from a gauge-invariant dimension-5 operator (LLHH) that is suppressed by a large UV cutoff scale  $\Lambda_\nu \sim v^2/m_\nu$ , far beyond the weak scale.
- However, this dimension-5 operator is nonrenormalizable and its Minimal UV Completion is given by the canonical seesaw with  $\Lambda_\nu = M_R$  after adding back  $v_R$  for each fermion family.
- The existence of  $v_R$  is predicted by the SM structure and provides the minimal UV completion for dimension-5 Weinberg operator through canonical Seesaw Mechanism, naturally generating light neutrino masses.
- Yet,  $v_R$  points to a brand-new seesaw scale  $\Lambda_\nu \sim v^2/m_\nu \sim 10^{14}\text{GeV}$  that is far beyond SM.
- →→ It is extremely important to probe  $v_R$  as the Last Missing Piece of the SM and test Neutrino Mass-Generation via canonical Seesaw Mechanism.
- SM prediction  $v_R$  could be wrong by EXP, but again its chance of success is high !

# Structures of 2 SM

## SM of Particle Physics + SM of Cosmology ( 粒子物理标准模型 + 宇宙学标准模型 )

### Standard Model of Elementary Particles

three generations of matter (fermions)			interactions / force carriers (bosons)	
mass charge spin	I U up	II C charm	III t top	0 0 1 g gluon
mass charge spin	I d down	II S strange	III b bottom	0 0 1 H higgs
LEPTONS	=4.7 MeV/c <sup>2</sup> -1/2 e electron	=95 MeV/c <sup>2</sup> -1/2 μ muon	=1.7768 GeV/c <sup>2</sup> -1/2 τ tau	0 0 1 γ photon
mass charge spin	<1.0 eV/c <sup>2</sup> 0 ν <sub>e</sub> electron neutrino	<0.17 MeV/c <sup>2</sup> 0 ν <sub>μ</sub> muon neutrino	<18.2 MeV/c <sup>2</sup> 0 ν <sub>τ</sub> tau neutrino	=124.97 GeV/c <sup>2</sup> 0 1 Z Z boson
				=80.433 GeV/c <sup>2</sup> ±1 W W boson
GAUGE BOSONS VECTOR BOSONS				
SCALAR BOSONS				

Higgs h(125GeV) (2012) is actually  
NOT the Last Building Block of SM .

But,  $\nu_R$  is the Last Missing Piece of SM !!

# GR - Einstein-Hilbert Action:

$$S = \int d^4x \sqrt{-g} [\kappa^{-2}(R + \Lambda) + \mathcal{L}_m]$$

## SM of Particle Physics + SM of Cosmology 粒子物理标准模型 + 宇宙学标准模型

### Standard Model of Elementary Particles

three generations of matter (fermions)			interactions / force carriers (bosons)	
I	II	III	g	H
mass =2.2 MeV/c <sup>2</sup> charge 2/3 spin 1/2 u up	mass =1.28 GeV/c <sup>2</sup> charge 2/3 spin 1/2 c charm	mass =173.1 GeV/c <sup>2</sup> charge 2/3 spin 1/2 t top	0 0 1 g gluon	0 0 1 H higgs
mass =4.7 MeV/c <sup>2</sup> charge -1/3 spin 1/2 d down	mass =96 MeV/c <sup>2</sup> charge -1/3 spin 1/2 s strange	mass =4.18 GeV/c <sup>2</sup> charge -1/3 spin 1/2 b bottom	0 0 1 γ photon	
mass =0.511 MeV/c <sup>2</sup> charge -1 spin 1/2 e electron	mass =105.66 MeV/c <sup>2</sup> charge -1 spin 1/2 μ muon	mass =1.7768 GeV/c <sup>2</sup> charge -1 spin 1/2 τ tau	0 1 Z Z boson	
mass <1.0 eV/c <sup>2</sup> charge 0 spin 1/2 ν <sub>e</sub> electron neutrino	mass <0.17 MeV/c <sup>2</sup> charge 0 spin 1/2 ν <sub>μ</sub> muon neutrino	mass <18.2 MeV/c <sup>2</sup> charge 0 spin 1/2 ν <sub>τ</sub> tau neutrino	0 1 W W boson	

SCALAR BOSONS

GAUGE BOSONS  
VECTOR BOSONS

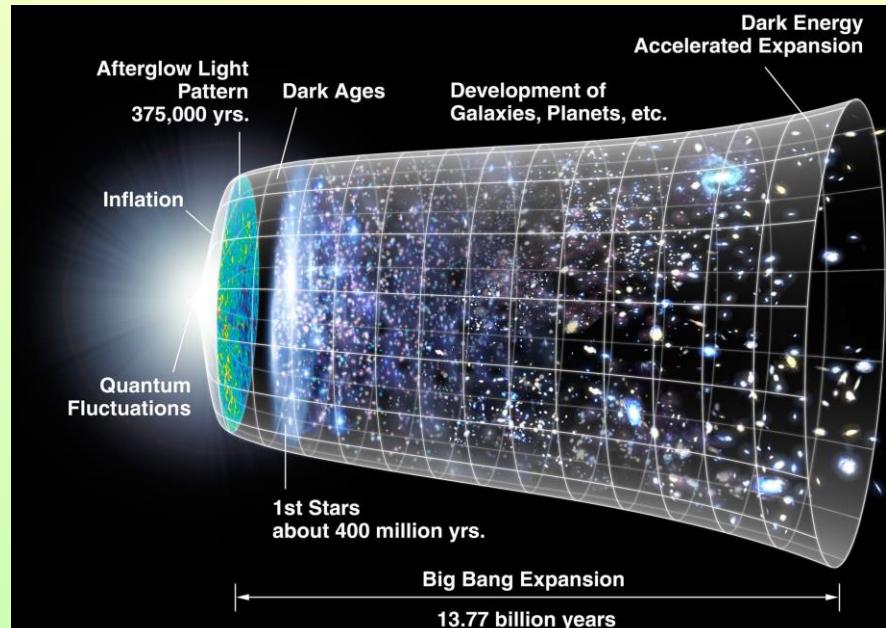
物质的基本组成 与 相互作用

Missing ν<sub>R</sub>

No DM??

No Dark Energy??

### Cosmological SM: ΛCDM + Inflation



物质的起源与演化

# Cosmological Bounds on light v Mass Sum

**Table 26.2:** Summary of  $\sum m_\nu$  constraints.

**PDG**

	Model	95% CL (eV)	Ref.
<b>CMB alone</b>			
Pl18[TT+lowE]	$\Lambda$ CDM+ $\sum m_\nu$	< 0.54	[24]
Pl18[TT,TE,EE+lowE]	$\Lambda$ CDM+ $\sum m_\nu$	< 0.26	[24]
<b>CMB + probes of background evolution</b>			
Pl18[TT,TE,EE+lowE] + BAO	$\Lambda$ CDM+ $\sum m_\nu$	< 0.13	[49]
Pl18[TT,TE,EE+lowE] + BAO	$\Lambda$ CDM+ $\sum m_\nu$ +5 params.	< 0.515	[25]
<b>CMB + LSS</b>			
Pl18[TT+lowE+lensing]	$\Lambda$ CDM+ $\sum m_\nu$	< 0.44	[24]
Pl18[TT,TE,EE+lowE+lensing]	$\Lambda$ CDM+ $\sum m_\nu$	< 0.24	[24]
Pl18[TT,TE,EE+lowE]+ ACT[lensing]	$\Lambda$ CDM+ $\sum m_\nu$	< 0.12	[50]
<b>CMB + probes of background evolution + LSS</b>			
Pl18[TT,TE,EE+lowE] + BAO + RSD	$\Lambda$ CDM+ $\sum m_\nu$	< 0.10	[49]
Pl18[TT,TE,EE+lowE+lensing] + BAO + RSD + Shape	$\Lambda$ CDM+ $\sum m_\nu$	< 0.082	[51]
Pl18[TT+lowE+lensing] + BAO + Lyman- $\alpha$	$\Lambda$ CDM+ $\sum m_\nu$	< 0.087	[52]
Pl18[TT,TE,EE+lowE] + BAO + RSD + SN + DES-Y1	$\Lambda$ CDM+ $\sum m_\nu$	< 0.12	[49]
Pl18[TT,TE,EE+lowE] + BAO + RSD + SN + DES-Y3	$\Lambda$ CDM+ $\sum m_\nu$	< 0.13	[53]



→ Finally, the **heaviest Light Neutrino Mass should be around 0.05eV (IO) — 0.06eV (NO).**

# Upper Bounds on Largest Light v Mass

$$\text{NO: } \sum m_\nu = m_3 + \sqrt{m_3^2 - \Delta m_{31}^2} + \sqrt{m_3^2 - \Delta m_{31}^2 + \Delta m_{21}^2}$$
$$\text{IO: } \sum m_\nu = m_2 + \sqrt{m_2^2 - \Delta m_{21}^2} + \sqrt{m_2^2 - |\Delta m_{32}|^2},$$

- **Largest Light Neutrino Mass:**
  - **Normal Ordering (NO): 0.06 eV**
  - **Inverted Ordering (IO): 0.05 eV**

# Beyond SM: Scale of Neutrino Mass Generation

- Weinberg Operator:  $1/\Lambda_\nu$  (LLHH)  
→→ Model-independent formulation of  $\nu$  Mass Generation !

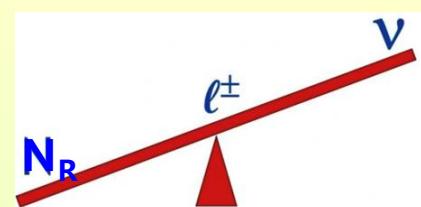
➤ Canonical Seesaw is the Minimal UV Completion.

➤ Seesaw Scale:  $\Lambda_\nu = M_R = m_D^2/m_\nu \sim 10^{14} \text{ GeV}$

→→ Seesaw Scale is a brand-new Scale beyond SM !

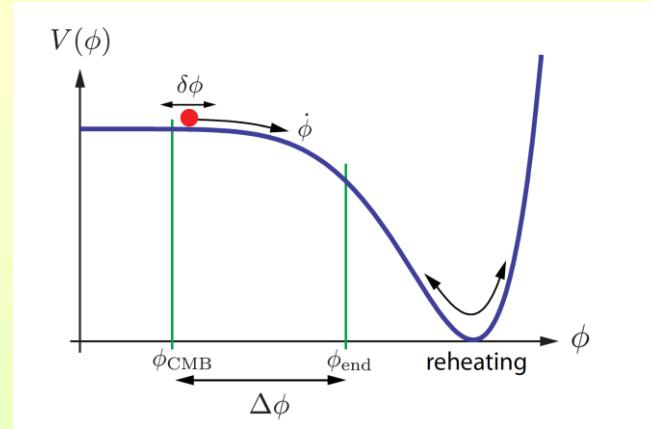
→ Great Challenge:  
How to test the High Scale Seesaw ? ?

## Neutrino Seesaw



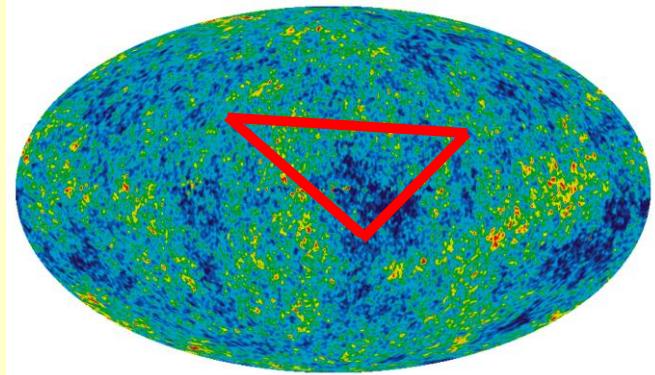
**Inflation + Seesaw occur around Same Scale  $\sim 10^{14}$  GeV !**

$$\Delta\mathcal{L} = \sqrt{-g} \left[ -\frac{1}{2} \partial_\mu \phi \partial^\mu - V(\phi) + \bar{N}_R i \bar{\partial} N_R + \frac{1}{\Lambda} \partial_\mu \phi \bar{N}_R \gamma^\mu \gamma^5 N_R \right. \\ \left. + \left( -\frac{1}{2} M \bar{N}_R^c N_R - y_\nu \bar{\ell}_L \tilde{H} N_R + \text{H.c.} \right) \right],$$



- $V(\phi)$  is the potential for inflation, is unknown but dominated by the mass term after inflation.
- Derivative coupling to keep the flatness of the inflaton potential (shift-symmetry).
- $\Lambda > 60 H_{\text{inf}}$  to keep perturbative unitarity.
- Inflaton coupling with SM fermions does not affect the analysis.
- After inflation, inflaton oscillates at the bottom of the potential until it decays into heavy  $N_R$  ( $M_\phi > 2 M_R$ ). The heavy neutrinos quickly decay into SM particles and reheat the universe.

# Non-Gaussianity



Non-Gaussianity is sensitive to New Physics !

- New physics could induce large non-Gaussianity: multi-field inflation models, modulated reheating, curvaton scenario .....
- Current limit from Planck on local type  $f_{NL} \sim \mathcal{O}(10)$ , future CMB observations, LiteBIRD  $\mathcal{O}(1)$ , large scale structure observations DESI  $\mathcal{O}(1)$ , 21cm tomography  $\mathcal{O}(0.01-0.1)$
- Non-Gaussianity can provide information to the new particle mass, spin, interactions: Cosmological Collider Signals

# Local type non-Gaussianity

- The local type non-Gaussianity which is defined by Bardeen Potential:

$$\Phi \equiv \frac{3}{5}\zeta$$

$$\langle \Phi_{\mathbf{k}_1} \Phi_{\mathbf{k}_2} \Phi_{\mathbf{k}_3} \rangle'_{\text{local}} = 2A^2 f_{\text{NL}}^{\text{local}} \left\{ \frac{1}{k_1^3 k_2^3} + \frac{1}{k_2^3 k_3^3} + \frac{1}{k_3^3 k_1^3} \right\}$$

- In the limit  $k_1 \sim k_2 \gg k_3$ , we find

$$f_{\text{NL}}^{\text{local}} \sim -\frac{10}{3} \frac{z_1^3 H^3}{(2\pi)^4 \mathcal{P}_\zeta^2} \cdot \left( \frac{\lambda \bar{h}}{2H} N_e - \frac{H \cdot z_2}{4z_1} \right)$$

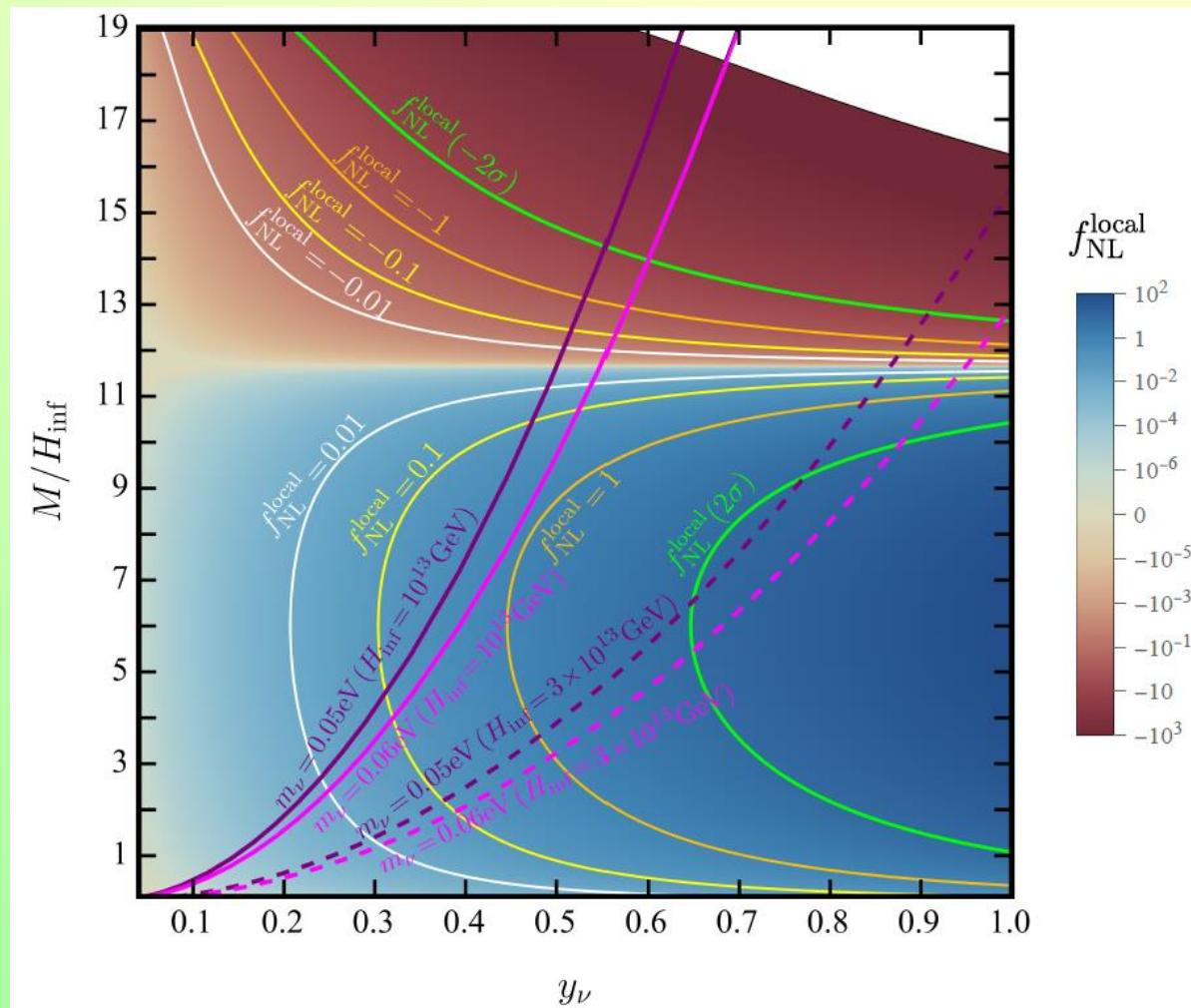
$$f_{\text{NL}}^{\text{local}} = -0.9 \pm 5.1 \quad (68\% \text{ C.L.}, \text{Planck 2018})$$

- Benchmark for this analysis:

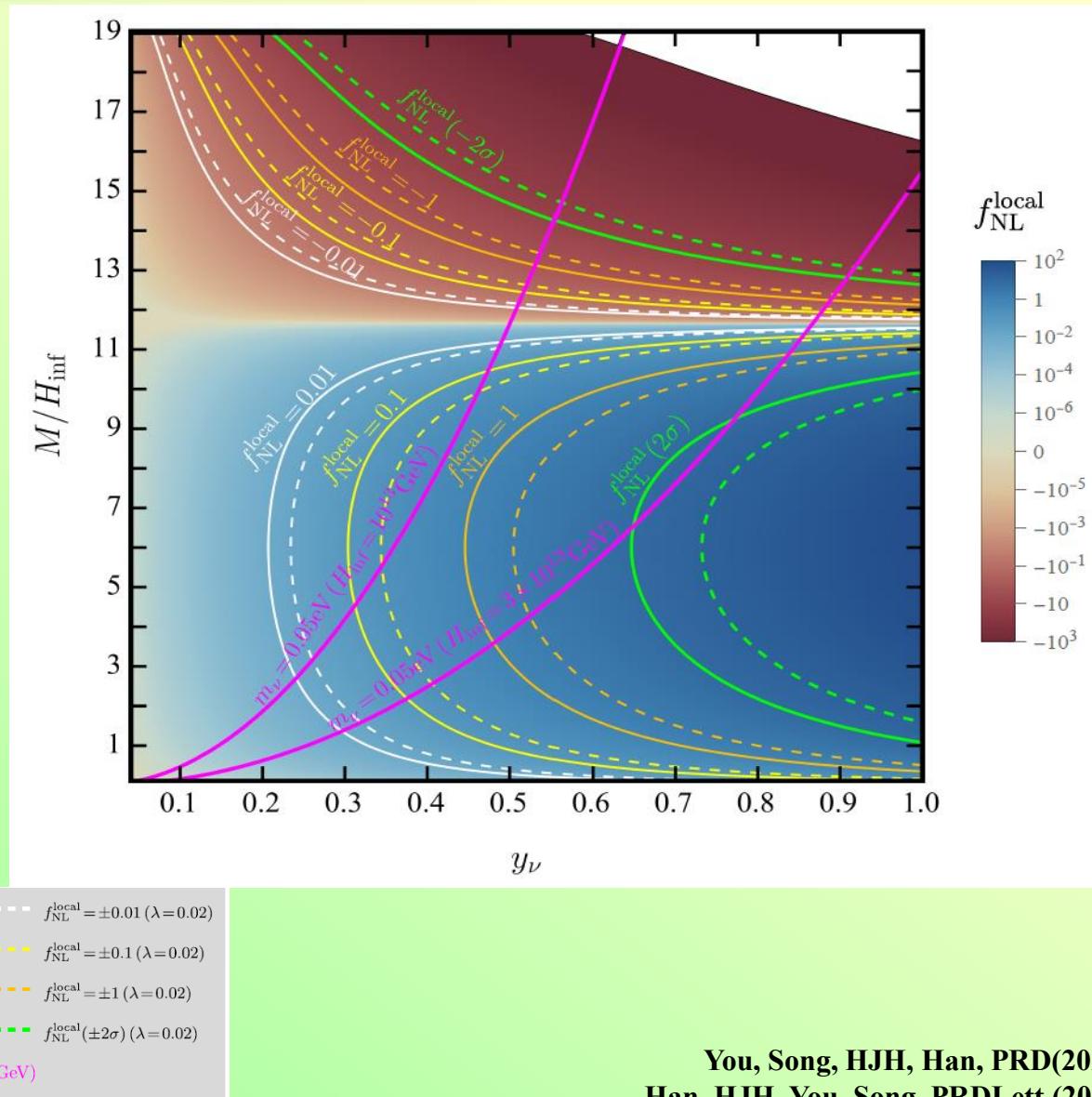
Parameters	$\mathcal{P}_\zeta$	$N_e$	$H_{\text{inf}}$	$m_\phi$	$\Lambda$	$\lambda$
Values	$2.1 \times 10^{-9}$	60	$(1, 3) \times 10^{13} \text{ GeV}$	$40H_{\text{inf}}$	$60H_{\text{inf}}$	0.01

# non-Gaussianity Probes Seesaw:

**$M_R$  vs Yukawa  $y_\nu$ ,  
and  $v_R$  vs  $v_L$  Mass Scales**



# NG: Seesaw & Higgs Self-Coupling



# Summary-1

- We propose a Minimal Model **incorporating inflation and seesaw** through  $\phi - v_R$  interaction.  
It provides a **New Realization of Higgs-modulated reheating**.
- It gives a **direct probe of Seesaw for Neutrino Mass Generation !**  
It links **Non-Gaussianity constraint to Low Energy v-Exps on light v masses** and to **LHC measurements on Higgs self-coupling**, via RG running.
- **Seesaw generated non-Gaussianity could be probed in near future CMB or large-scale structure observations**
- It also provides a framework of cosmological Higgs collider (**particles coupling to Higgs boson could be detected**)

# Geometric Higgs Mechanism for KK Graviton Mass Generation

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Hang and HJH, PRD(2022), 2106.04568

Research (2022), 2207.11214

Hang, Zhao, HJH, Qiu, (2024), 2406.12713

# Challenges of Mass Generations in 4d

Nothing forbids **Spin-0 Higgs Boson** gets a huge mass from quantum corrections!  
**Quark/Lepton mass-scales** are protected by P-violation and set by Higgs VEV through Yukawa coupling.

Fields	$m = 0$	$m \neq 0$
$H^0$ ( $s=0$ )	1 ( $\lambda = 0$ )	= ( $\lambda = 0$ )
$A^\mu$ ( $s=1$ )	2 ( $\lambda = \pm 1$ )	$\neq$ ( $\lambda = 0, \pm 1$ )
$h^{\mu\nu}$ ( $s=2$ )	2 ( $\lambda = \pm 2$ )	$\neq$ ( $\lambda = 0, \pm 1, \pm 2$ )
$(\psi_L, \psi_R)$ ( $s=\frac{1}{2}$ )	$1 + 1$ ( $\lambda = \frac{1}{2}, -\frac{1}{2}$ )	2 ( $\lambda = \pm \frac{1}{2}$ )

- “Higgs” Mechanism can be
  - 1). Conventional, 2). Geometric, 3). Topological

-- Conventional Higgs Mechanism:

Mass-Generation by SSB by Vacuum Expectation Value of Higgs Boson.

-- Geometric “Higgs” Mechanism:

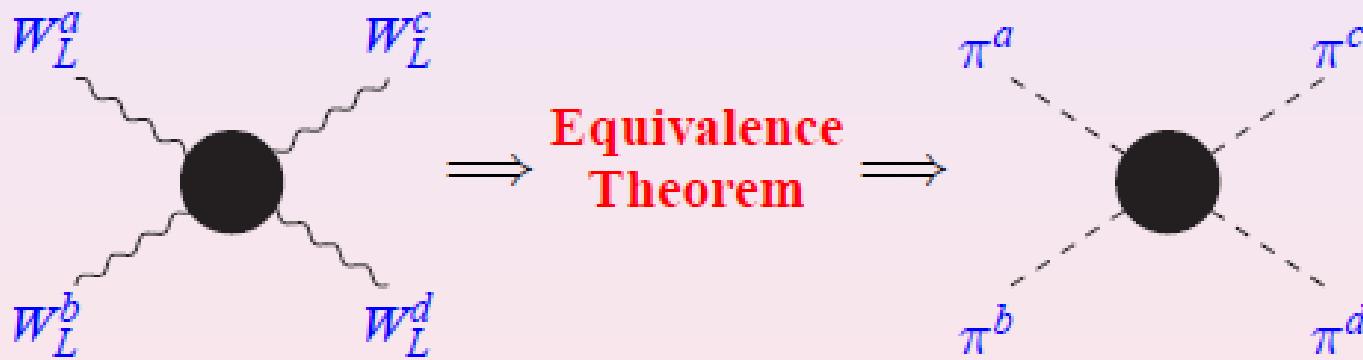
Mass-Generation SSB by Compactification of Extra Dimension.

-- Topological “Higgs” Mechanism:

No SSB, construction of Gauge-Invariant Mass Term with topology,  
Dynamical Conversion of Physical Degrees of Freedom.

## No Lose Theorem for LHC Discovery

— Equivalence Theorem —  
the Bridge between Experiments and EWSB



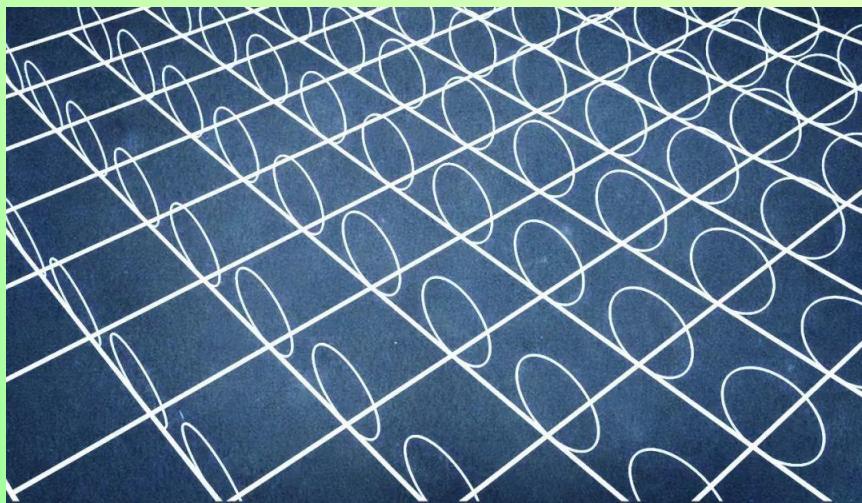
Energy Cancellations:  $E^4 \rightarrow E^0$

Unitarity Bound  $\rightarrow m_H < 800\text{GeV}$  or  $\Lambda < 1.2\text{TeV}$

For a review: H.-J. He et al, arXiv:hep-ph/9704276.

## 5d: Kaluza-Klein Compactification

- We are apparently living in (3+1)d spacetime.
- But our Universe could have **Extra Dimensions** beyond  $d = 4$  ,  
except that **Extra Dimensions are all curled-up !**
- Simplest case is a **5<sup>th</sup> dimension** curled up on a **circle  $S^1$** ,  
called **Kaluza-Klein (KK) Compactification** (in 1920s)



# Geometric “Higgs” Mechanism in 5d

➤ Under compactification of flat 5d, how do **KK Masses** arise? ---

► Consider 5d Massless Gauge Fields  $\hat{A}^{aM}(x^\mu, x^5)$ ,

$$\left\{ \begin{array}{l} 0 = P^2 = p_\mu p^\mu + p_5 p^5 = p^2 - p_5^2 \\ p^5 = \frac{n}{R}, \quad (n = 0, 1, 2, \dots), \end{array} \right. \quad \text{(After 5d compactification)}$$


where  $p^5$  is quantized due to BCs, and  $p^2 = p_\mu p^\mu$  is the 4-momentum-squared in 4d.

► Hence, in 4d we can see the “**KK Tower**”, as a unique consequence of **5d Compactification**,

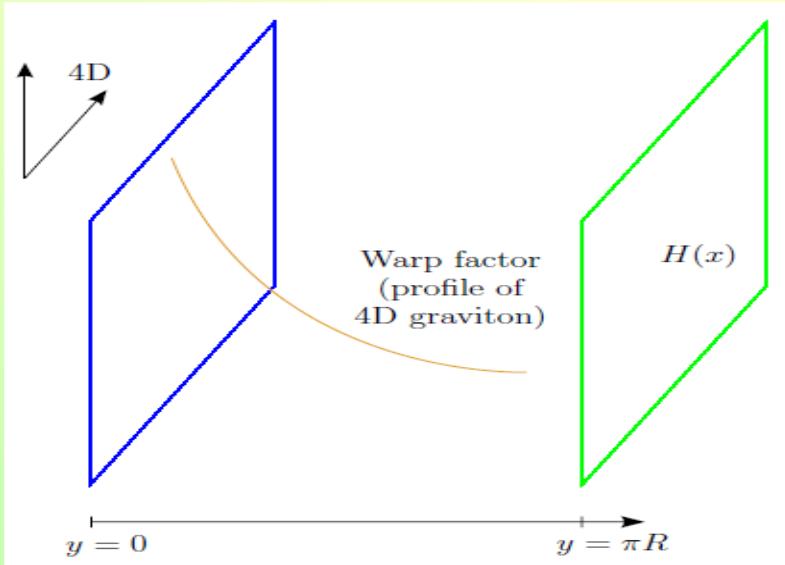
$$p^2 = p_5^2 = \frac{n^2}{R^2}, \quad (n = 0, 1, 2, \dots), \quad \Rightarrow \quad \text{KK Tower!}$$

➤ Masses Generation by KK Compactification:

—————> Geometric “Higgs” Mechanism

# Warped 5d: Geometric “Higgs” Mechanism

## ➤ Warped 5d RS1 under $S^1/Z_2$ :



Randall-Sundrum  
PRL.1999, hep-ph/9905221

## ➤ Warped 5d metric:

$$\begin{aligned} ds^2 &= e^{2A(y)} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2, \\ ds^2 &= e^{2A(z)} (\eta_{\mu\nu} dx^\mu dx^\nu + dz^2) \end{aligned}$$

where

$$A(z) = -ky, \quad A(z) = -\ln(1 + kz), \quad k = \sqrt{-\Lambda/6} \quad \hat{g}_{MN} = e^{2A(z)} \hat{\eta}_{MN}$$

# Geometric Higgs Mechanism for KK Gravitons

- Under  $S^1/Z_2$  compactification, we impose the Boundary Conditions:

$$\partial_5 \hat{h}_{\mu\nu} \Big|_{x^5=0,L} = 0, \quad \partial_5 \hat{\phi} \Big|_{x^5=0,L} = 0, \quad \hat{h}_{\mu 5} \Big|_{x^5=0,L} = 0$$

- 4d Quadratic KK Lagrangian contains Mixing Terms:

$$+ 2\mathbb{M}_n h_n \partial_\mu \mathcal{V}_n^\mu - 2\mathbb{M}_n h_n^{\mu\nu} \partial_\mu \mathcal{V}_{\nu,n} - \frac{3}{2} \mathbb{M}_n^2 h_n \phi_n - 3\mathbb{M}_n \partial_\mu \mathcal{V}_n^\mu \phi_n$$

which can be eliminated by  $R_\xi$  gauge fixing:

$$\mathcal{V}_n^\mu = h_n^{\mu 5}$$

$$\phi_n = h_n^{55}$$

→  $\mathcal{L}_{\text{GF}} = \int_0^L dz \hat{\mathcal{L}}_{\text{GF}} = - \sum_{n=0}^{\infty} \frac{1}{\xi_n} \left[ (\mathcal{F}_n^\mu)^2 + (\mathcal{F}_n^5)^2 \right],$

→  $\mathcal{F}_n^\mu = \partial_\nu h_n^{\mu\nu} - \left( 1 - \frac{1}{2\xi_n} \right) \partial^\mu h_n + \frac{1}{\sqrt{2}} \xi_n \mathbb{M}_n \mathcal{V}_n^\mu,$   
 $\mathcal{F}_n^5 = \frac{1}{2} \mathbb{M}_n h_n - \sqrt{\frac{3}{2}} \xi_n \mathbb{M}_n \phi_n + \frac{1}{\sqrt{2}} \partial_\mu \mathcal{V}_n^\mu.$

➤ R<sub>ξ</sub> Gauge Propagators:

$$\begin{aligned}
 \mathcal{D}_{nm}^{\mu\nu\alpha\beta}(p) = & -\frac{i\delta_{nm}}{2} \left\{ \frac{(\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta})}{p^2 + M_n^2} \right. \\
 & + \frac{1}{3} \left[ \frac{1}{p^2 + M_n^2} - \frac{1}{p^2 + (3\xi_n - 2)M_n^2} \right] \left( \eta^{\mu\nu} - \frac{2p^\mu p^\nu}{M_n^2} \right) \left( \eta^{\alpha\beta} - \frac{2p^\alpha p^\beta}{M_n^2} \right) \\
 & + \frac{1}{M_n^2} \left[ \frac{1}{p^2 + M_n^2} - \frac{1}{p^2 + \xi_n M_n^2} \right] (\eta^{\mu\alpha}p^\nu p^\beta + \eta^{\mu\beta}p^\nu p^\alpha + \eta^{\nu\alpha}p^\mu p^\beta + \eta^{\nu\beta}p^\mu p^\alpha) \\
 & \left. + \frac{4p^\mu p^\nu p^\alpha p^\beta}{\xi_n M_n^4} \left( \frac{1}{p^2 + \xi_n^2 M_n^2} - \frac{1}{p^2 + \xi_n M_n^2} \right) \right\}, \tag{2.1}
 \end{aligned}$$

# KK Graviton Propagator in Feynman/Unitary Gauge

➤ Feynman Gauge ( $\xi_n=1$ ) Propagators take very simple form:

$$\mathcal{D}_{nm}^{\mu\nu\alpha\beta}(p) = -\frac{i\delta_{nm}}{2} \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2 + M_n^2}$$

➤ We also derive Unitary Gauge ( $\xi_n=\infty$ ) Propagator:

$$\mathcal{D}_{nm,UG}^{\mu\nu\alpha\beta}(p) = -\frac{i\delta_{nm}}{2} \frac{\bar{\eta}^{\mu\alpha}\bar{\eta}^{\nu\beta} + \bar{\eta}^{\mu\beta}\bar{\eta}^{\nu\alpha} - \frac{2}{3}\bar{\eta}^{\mu\nu}\bar{\eta}^{\alpha\beta}}{p^2 + M_n^2}$$

where  $\bar{\eta}^{\mu\nu} = \eta^{\mu\nu} + p^\mu p^\nu / M_n^2$ .

# KK Graviton Propagator vs vDVZ Discontinuity

- KK Graviton propagator in  $M_n \rightarrow 0$  limit:

$$\begin{aligned}
 \mathcal{D}_{nm}^{\mu\nu\alpha\beta}(p) &= -\frac{i\delta_{nm}}{2} \left[ \frac{(\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta})}{p^2} - \frac{1-\xi_n}{p^4} (\eta^{\mu\alpha}p^\nu p^\beta + \eta^{\mu\beta}p^\nu p^\alpha + \eta^{\nu\alpha}p^\mu p^\beta \right. \\
 &\quad \left. + \eta^{\nu\beta}p^\mu p^\alpha - 2\eta^{\mu\nu}p^\alpha p^\beta - 2\eta^{\alpha\beta}p^\mu p^\nu) - 4(1-\xi_n)^3 \frac{p^\mu p^\nu p^\alpha p^\beta}{p^6} \right], \\
 &= -\frac{i\delta_{nm}}{2} \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2}, \quad (\text{for } \xi_n = 1).
 \end{aligned}$$



→→ It is **free from vDVZ (van-Dam-Veltman-Zakharov) Discontinuity !**

- Compared with **conventional massless graviton propagator:**

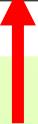
$$\begin{aligned}
 \mathcal{D}_{00}^{\mu\nu\alpha\beta}(p) &= -\frac{i}{2} \left[ \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2} + (\xi - 1) \frac{\eta^{\mu\alpha}p^\nu p^\beta + \eta^{\mu\beta}p^\nu p^\alpha + \eta^{\nu\alpha}p^\mu p^\beta + \eta^{\nu\beta}p^\mu p^\alpha}{p^4} \right] \\
 &= -\frac{i}{2} \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2}, \quad (\text{for } \xi = 1).
 \end{aligned}$$



# KK Graviton Propagator vs vDVZ Discontinuity

- KK Graviton propagator in  $M_n \rightarrow 0$  limit:

$$\begin{aligned}
 D_{nm}^{\mu\nu\alpha\beta}(p) &= -\frac{i\delta_{nm}}{2} \left[ \frac{(\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta})}{p^2} - \frac{1-\xi_n}{p^4} (\eta^{\mu\alpha}p^\nu p^\beta + \eta^{\mu\beta}p^\nu p^\alpha + \eta^{\nu\alpha}p^\mu p^\beta \right. \\
 &\quad \left. + \eta^{\nu\beta}p^\mu p^\alpha - 2\eta^{\mu\nu}p^\alpha p^\beta - 2\eta^{\alpha\beta}p^\mu p^\nu) - 4(1-\xi_n)^3 \frac{p^\mu p^\nu p^\alpha p^\beta}{p^6} \right], \\
 &= -\frac{i\delta_{nm}}{2} \frac{\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha} - \eta^{\mu\nu}\eta^{\alpha\beta}}{p^2}, \quad (\text{for } \xi_n = 1).
 \end{aligned}$$



- Compared with Fierz-Pauli massive graviton propagator:

$$D_{\text{PF}}^{\mu\nu\alpha\beta}(p) = -\frac{i}{2} \frac{\bar{\eta}^{\mu\alpha}\bar{\eta}^{\nu\beta} + \bar{\eta}^{\mu\beta}\bar{\eta}^{\nu\alpha} - \frac{2}{3}\bar{\eta}^{\mu\nu}\bar{\eta}^{\alpha\beta}}{p^2 + M^2}$$

where  $\bar{\eta}^{\mu\nu} = \eta^{\mu\nu} + p^\mu p^\nu/M^2$ .

It has vDVZ discontinuity in  $M_n \rightarrow 0$  limit!

# KK Graviton Propagator vs vDVZ Discontinuity

➤ 5d KK Graviton conserves physical d.o.f in  $M_n \rightarrow 0$  limit:

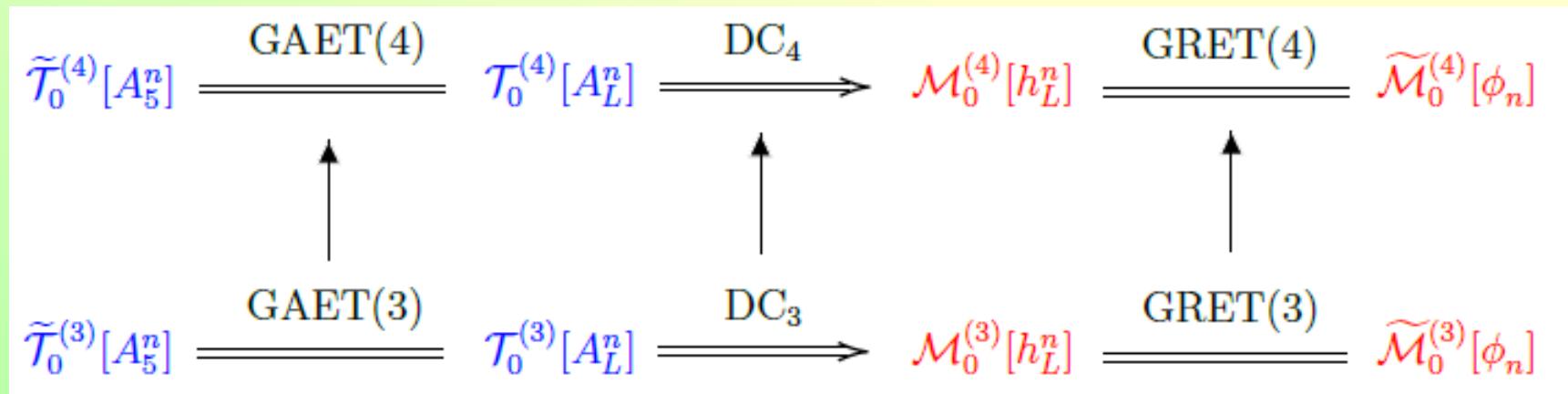
$$5 = 2 + 2 + 1$$
$$\lambda = \begin{matrix} 0, \pm 1, \pm 2 \\ \pm 2 \\ \pm 1 \\ 0 \end{matrix}$$

where  $\lambda = 0, \pm 1$  d.o.f arise from 3 KK Goldstones  $(h_n^{\mu 5}, h_n^{55})$ .

➤ But, the Massive Graviton of Fierz-Pauli gravity does not conserve physical d.o.f in  $M_n \rightarrow 0$  limit:

$$5 \neq 2$$

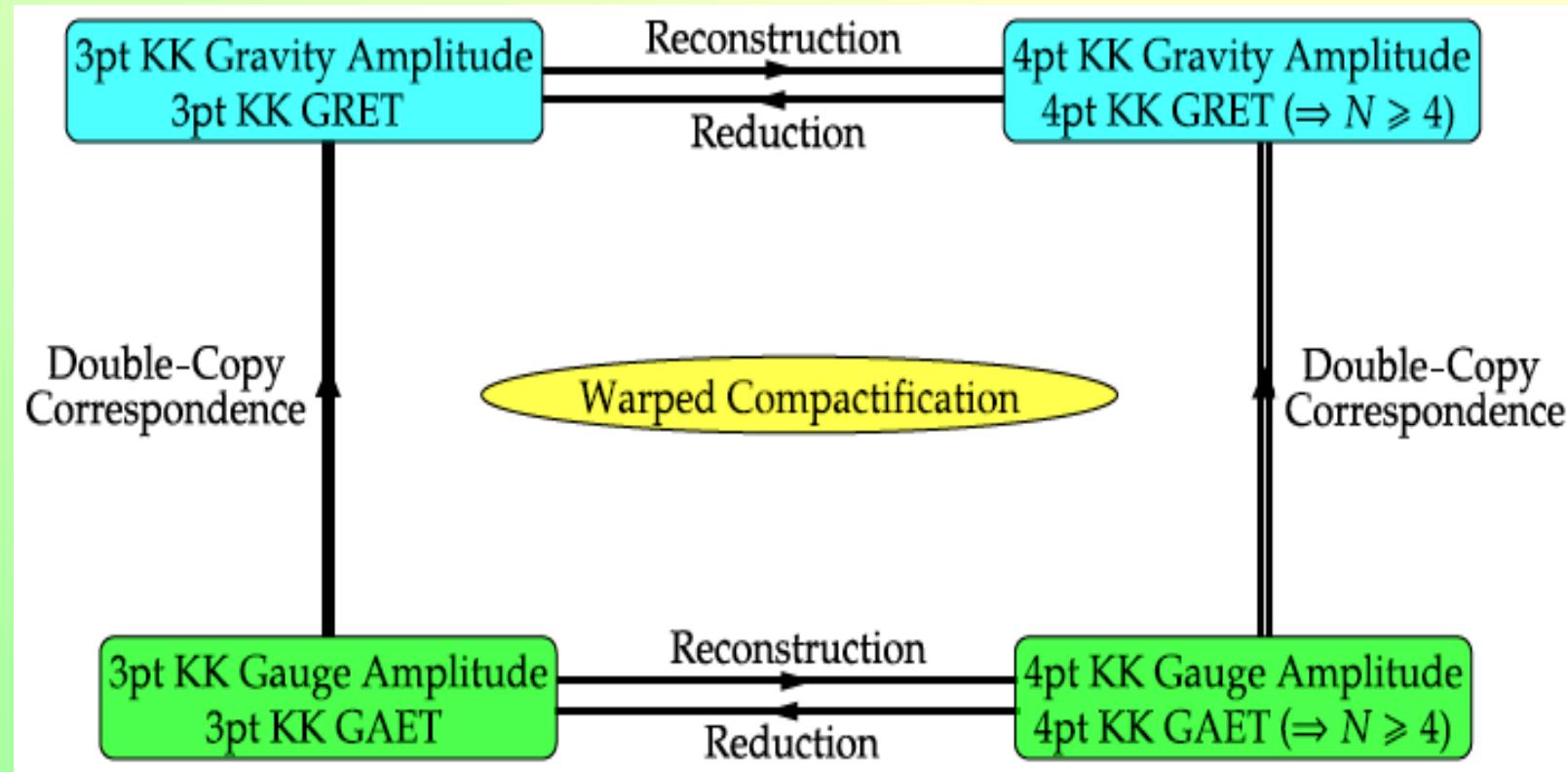
# Schematic Summary: ET vs Double-Copy and from 3pt to 4pt Amplitudes



- This can be further extended to the case of  $N > 4$  amplitudes.

# KK ET & Double-Copy Correspondance: 3-Point to 4-Point

- Equivalence Theorem and Double-Copy Correspondance  
from 3pt KK amplitudes to 4pt KK amplitudes and  
from massive KK gauge amplitudes to massive KK gravitational amplitudes.



# Geometric Higgs Mechanism via GRET

➤ **Gravitational Equivalence Theorem (GRET type-I):**

$$\mathcal{M}[h_{n_1}^L, \dots, h_{n_N}^L, \Phi] = C_{\text{mod}}^{n_j m_j} \mathcal{M}[\phi_{m_1}, \dots, \phi_{m_N}, \Phi] + O(\mathbb{M}_n/E_n - \text{suppressed})$$

$$C_{\text{mod}}^{n_j m_j} = C_{n_1 m_1} \cdots C_{n_N m_N} = \delta_{n_1 m_1} \cdots \delta_{n_N m_N} + O(\text{loop})$$

➤ **Gravitational Equivalence Theorem (GRET type-II):**

$$\mathcal{M}[h_{n_1}^{\pm 1}, \dots, h_{n_N}^{\pm 1}, \Phi] = \hat{C}_{\text{mod}}^{n_j m_j} \mathcal{M}[\mathcal{V}_{m_1}^{\pm 1}, \dots, \mathcal{V}_{m_N}^{\pm 1}, \Phi] + O(\mathbb{M}_n/E_n - \text{suppressed})$$

$$\hat{C}_{\text{mod}}^{n_j m_j} = \hat{C}_{n_1 m_1} \cdots \hat{C}_{n_N m_N} = (-i)^N \delta_{n_1 m_1} \cdots \delta_{n_N m_N} + O(\text{loop})$$

# Generalized Power Counting Rule for KK Theories

## ➤ Generalized Energy Power Counting for KK Theories:

$$D_E = 2\mathcal{E}_{h_L} + (2L+2) + \sum_j \mathcal{V}_j (d_j - 2 + \frac{1}{2}f_j).$$

## ➤ E-Power Counting of helicity-0 KK Graviton/Goldstone Amplitudes:

$$\begin{aligned} D_E(Nh_n^L) &= 2(N+L+1), & D_E(N\phi_n) &= 2(L+1), \\ \implies D_E(Nh_n^L) - D_E(N\phi_n) &= 2N, \end{aligned} \tag{1}$$

## ➤ E-Power Counting of helicity-1 KK Graviton/Goldstone Amplitudes:

$$\begin{aligned} D_E(Nh_n^{\pm 1}) &= N+2(L+1), & D_E(N\mathcal{V}_n^{\pm 1}) &= 2(L+1), \\ \implies D_E(Nh_n^{\pm 1}) - D_E(N\mathcal{V}_n^{\pm 1}) &= N. \end{aligned}$$

## ➤ E-Power Counting of KK Gauge/Goldstone boson Amplitudes:

$$\begin{aligned} D_E(NA_L^{an}) &= 4, & D_E(NA_5^{an}) &= 4 - N - \bar{V}_3^{\min}, \\ \implies D_E(NA_L^{an}) - D_E(NA_5^{an}) &= N + \bar{V}_3^{\min} \end{aligned}$$

- **Gravitational Equivalence Theorem (GET) identity:**

$$\mathcal{M}[h_{n_1}^L(k_1), \dots, h_{n_N}^L(k_N), \Phi] = \mathcal{M}[\phi_{n_1}(k_1), \dots, \phi_{n_N}(k_N), \Phi] + \sum_{1 \leq j \leq N} \mathcal{M}[\{\tilde{\Delta}_{n_j}, \phi_n\}, \Phi],$$

- **Energy Power Counting:**  $D_E[Nh_n^L] = 2(N+1) + 2L,$

$$D_E[N\phi_n] = 2 + 2L. \quad D_E[N\tilde{v}_n] = 2 + 2L,$$

- **We deduce  $h_L$  Amplitude has Large E-Cancellations:**

—————→  $D_E[Nh_n^L] - D_E[N\phi_n] = 2N$

- **For  $N = 4$  scattering at tree level, GET proves Large E-cancellation:**

—————→  $E^{10} \rightarrow E^2 \quad (\text{by } 10 - 2 = 8 \text{ powers})$

# KK Gauge Amplitudes & E-Cancellations at 3-Point

- 3pt Longitudinal KK gauge boson amplitude ( $A_L A_L A_T$ ) and KK Goldstone amplitude ( $A_5 A_5 A_T$ ):

$$\mathcal{T}_0[A_L^{an_1} A_L^{bn_2} A_{\pm}^{cn_3}] = -ig f^{abc}(\epsilon_3 \cdot p_1) \frac{(M_{n_3}^2 - M_{n_1}^2 - M_{n_2}^2)}{M_{n_1} M_{n_2}} a_{n_1 n_2 n_3},$$

$$\tilde{\mathcal{T}}_0[A_5^{an_1} A_5^{bn_2} A_{\pm}^{cn_3}] = -i2g f^{abc}(\epsilon_3 \cdot p_1) \tilde{a}_{n_1 n_2 n_3}$$

- The most fundamental GAET:



$$\mathcal{T}_0[A_L^{an_1} A_L^{bn_2} A_{\pm}^{cn_3}] = -\tilde{\mathcal{T}}_0[A_5^{an_1} A_5^{bn_2} A_{\pm}^{cn_3}]$$

enforcing energy cancellation:  $E^3 \rightarrow E^1$

and it requires the coupling-mass condition:

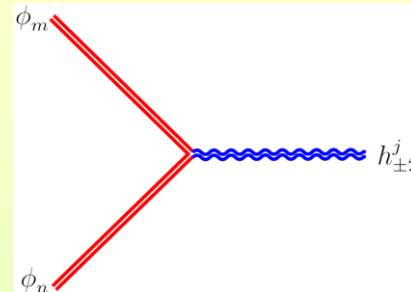
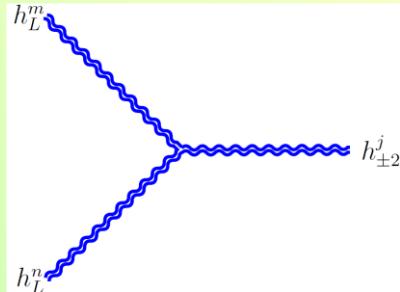


$$(M_{n_1}^2 + M_{n_2}^2 - M_{n_3}^2) a_{n_1 n_2 n_3} = 2M_{n_1} M_{n_2} \tilde{a}_{n_1 n_2 n_3}$$

from which all N-point KK Amplitudes and GAET (N>3) can be derived !

# KK Graviton Amplitudes & E-Cancellations at 3-Point

➤ 3pt Longitudinal KK graviton amplitude ( $h_L h_L h_T$ ) and KK Goldstone amplitude ( $\phi_n \phi_m h_T$ ) :



$$\mathcal{M}_0[h_{n_1}^L h_{n_2}^L h_{n_3}^{\pm 2}] = \frac{\kappa(\epsilon_3 \cdot p_1)^2}{6 \mathbb{M}_{n_1}^2 \mathbb{M}_{n_2}^2} \left[ 2 \mathbb{M}_{n_1}^2 \mathbb{M}_{n_2}^2 + (\mathbb{M}_{n_1}^2 + \mathbb{M}_{n_2}^2 - \mathbb{M}_{n_3}^2)^2 \right] \alpha_{n_1 n_2 n_3},$$

$$\tilde{\mathcal{M}}_0[\phi_{n_1} \phi_{n_2} h_{n_3}^{\pm 2}] = \kappa(\epsilon_3 \cdot p_1)^2 \tilde{\beta}_{n_1 n_2 n_3},$$

➤ E Cancellations for above 3-point KK Graviton Amplitude:  $E^6 \rightarrow E^2$

# Gauge-Gravity Duality: Double-Copy Massless vs Massive

Li, Hang, HJH, JHEP (2022), 2209.11191  
Hang, Zhao, HJH, Qiu, JHEP (2025), 2406.12713  
Hang, HJH, PRD (2022), 2106.04568, 2207.11214  
Li, Hang, HJH, He, JHEP(2022), 2111.12042  
Hang, HJH, Shen, JHEP(2022), 2110.05399  
Research (2023), 2406.13671  
H.X. Liu, Z.X. Yi, HJH, **2512.10870 [hep-th]** (66pp)

$$(\text{GR}) = (\text{QCD})^2$$

$$(\text{Gravity}) = (\text{Gauge Theory})^2$$

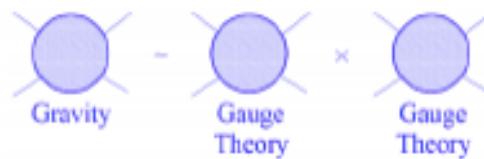
$$(\text{引力}) = (\text{规范力})^2$$

# Double-Copy: KLT vs BCJ

- for Amplitudes of **Massless open/closed strings** and
- for **massless gauge/graviton bosons**:

1985: Kawai, Lewellen, Tye (KLT): “closed string amp=open-string amp^2”

Field-theory limit:



$$M_4^{\text{tree}}(1,2,3,4) = -i g_{12} A_4^{\text{tree}}(1,2,3,4) A_4^{\text{tree}}(1,2,4,3)$$

gravity

gauge theory color ordered

Yang-Mills  
gauge theory:



Einstein  
gravity:



$$= g f^{abc} (\eta_{\mu\nu} (k_1 - k_2)_\rho + \text{cyclic})$$

$$i\kappa (\eta_{\mu\nu} (k_1 - k_2)_\rho + \text{cyclic}) \times (\eta_{\alpha\beta} (k_1 - k_2)_\gamma + \text{cyclic})$$

“square” of  
Yang-Mills  
vertex.

2008: Bern, Carrasco, Johansson (BCJ):

$$\mathcal{A}_4^{\text{tree}} = g^2 \left( \frac{n_s \mathcal{C}_s}{s} + \frac{n_t \mathcal{C}_t}{t} + \frac{n_u \mathcal{C}_u}{u} \right)$$

$$\mathcal{A}_4^{\text{tree}} \Big|_{c_i \rightarrow n_i} \equiv -i \mathcal{M}_4^{\text{tree}} = \frac{n_s^2}{s} + \frac{n_t^2}{t} + \frac{n_u^2}{u}$$



$$\begin{aligned} \mathcal{C}_s + \mathcal{C}_t + \mathcal{C}_u &= 0 \\ n_s + n_t + n_u &= 0 \end{aligned}$$

# Massive Double-Copy vs Mass Spectral Condition

- For the massive KK graviton amplitude:

$$\mathcal{M} \left[ h_{\sigma_1}^{n_1} h_{\sigma_2}^{n_2} h_{\sigma_3}^{n_3} h_{\sigma_4}^{n_4} \right] = -\frac{\kappa^2}{16} \sum_j \sum_{\lambda_k, \lambda'_k} \left( \prod_{k=1}^4 C_{\lambda_k \lambda'_k}^{\sigma_k} \right) \frac{\mathcal{N}_j^P(\lambda_k) \mathcal{N}_j^P(\lambda'_k)}{s_j - M_{nn_j}^2},$$

we impose the Generalized Gauge Transformation:

$$\mathcal{N}_j^P' = \mathcal{N}_j^P + (s_j - M_{nn_j}^2) \times \Delta,$$

from which we deduce conditions:

$$\sum_j \mathcal{N}_j^P = 0, \quad \sum_j (s_j - M_{nn_j}^2) = 0.$$

- This leads to the 4-point Mass Spectral Condition: → Nontrivial !  
→ Does Not always hold !

$$\sum_{i=1}^4 M_{n_i}^2 = M_{nn_s}^2 + M_{nn_t}^2 + M_{nn_u}^2$$

# Massive Double-Copy vs Mass Spectral Condition

- Start from a general 4-point **Mass Spectral Condition**:

$$M_1^2 + M_2^2 + M_3^2 + M_4^2 = M_{12}^2 + M_{13}^2 + M_{14}^2$$

- We ask: **How to solve it and What is the solution ??**
- For the **flat 5d Toroidal Compactification** of  $S^1$ , we have:

$$n_1^2 + n_2^2 + n_3^2 + n_4^2 = (n_1 + n_2)^2 + (n_1 + n_3)^2 + (n_1 + n_4)^2$$

which leads to the condition:

$$n_1 + n_2 + n_3 + n_4 = 0$$

- This is just the **KK number (5d momentum) Conservation!**
- E.g., it does not hold for **5d Orbifold Compactification** or **5d Warped Compactification** ! (additional treatments needed.)

# Massive Double-Copy vs Mass Spectral Condition

➤ Start from a general 4-point Mass Spectral Condition:

$$M_1^2 + M_2^2 + M_3^2 + M_4^2 = M_{12}^2 + M_{13}^2 + M_{14}^2$$

➤ We ask: How to solve it and What is the solution ???

➤ In fact, starting with 3 rather modest conditions:

- 1). A massive theory contains at least 2 types of particles with Unequal masses.
- 2). There exists only a simple pole in each of (s, t, u) channels.
- 3). Each scattering amplitude should include the contributions from all 3 kinematic channels of (s, t, u).

we can prove:

→→ Toroidal Compactification of flat Extra dimensions is the Unique Solution !

# Massive Double-Copy vs Mass Spectral Condition

- General 4-point Mass Spectral Condition:

$$M_1^2 + M_2^2 + M_3^2 + M_4^2 = M_{12}^2 + M_{13}^2 + M_{14}^2 \quad (1)$$

- Proof of the Unique Solution:

We identified the group structure underlying condition (1) is a product of Integer Groups  $\mathbb{Z}^r$  (with rank  $r$ ) in the Finitely Generated Abelian Group:

$$\mathcal{G} \cong \mathbb{Z}^r \oplus \mathbb{Z}_{p_1} \oplus \mathbb{Z}_{p_2} \oplus \cdots \oplus \mathbb{Z}_{p_s}, \quad (r, s \in \mathbb{N})$$

- Then, we proved that the unique solution to Condition (1) is given by  $M_{\{k\}}^2 = k^2 M_{\{1\}}^2$ , for the case  $r = 1$  and by  $M_n^2 = n \bar{M}^2 n^T$  for the case  $r > 1$ .
- By inspecting all the known consistent QFTs, we conclude that only the KK theories with  $\delta$  ( $= r$ ) extra dimensions under the toroidal compactification could have a mass spectrum behave exactly as above. These theories also hold the KK number conservation.

# Gauge Forces vs Gravity Force

- **Massive Case:**  $GR = (Gauge\ Theory)^2$  ??
- What happens to massive Kaluza-Klein Theory ??
- **First Principle Approach:** Using KK bosonic string theory, we derived Massive KLT Relations between product of KK Open String Amplitudes and KK Closed String Amplitude. Taking Field Theory Limit, we derived Massive KLT Relations between product of KK Gauge Boson Amplitudes KK Graviton Amplitude:

$$\mathcal{M}[1_L^n 2_L^n 3_L^n 4_L^n] = \frac{\kappa^2}{32} \sum_{\{a_j, b_j\}} \hat{\varrho}_{ab} \left\{ (s - 4M_n^2) \mathcal{T}_{a_j}[1^{+n} 2^{+n} 3^{-n} 4^{-n}] \mathcal{T}_{b_j}[1^{+n} 2^{+n} 4^{-n} 3^{-n}] \right. \\ \left. + s \mathcal{T}_{a_j}[1^{+n} 2^{-n} 3^{+n} 4^{-n}] \mathcal{T}_{b_j}[1^{+n} 2^{-n} 4^{+n} 3^{-n}] \right. \\ \left. + s \mathcal{T}_{a_j}[1^{+n} 2^{-n} 3^{-n} 4^{+n}] \mathcal{T}_{b_j}[1^{+n} 2^{-n} 4^{-n} 3^{+n}] \right\}$$

# Topological Mass Generation & Scattering Amplitudes Topological Equivalence Theorem & Double-Copy for Chern-Simons Gauge & Gravity Theories

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H.X. Liu, Z.X. Yi, HJH, [arXiv:2512.10870 \[hep-th\]](https://arxiv.org/abs/2512.10870) (66 pp)  
Hang, HJH, Shen JHEP(2022), 2110.05399  
Research (2023), 2406.13671

# 3d Topological Massive Gauge & Gravity Theories

- 3d Spacetime has distinctive features characterized by **Gauge & Gravitational Chern-Simons terms**, which generate **topological masses** for gauge bosons and gravitons, realize **fractional statistics** and predict the existence of **anyon-like** quasi-particles.
- Chern-Simons (CS) terms are topological invariants in mathematics and serve as the theoretical key ingredients of a wide range of applications in modern physics, including fractional quantum Hall effect, models of high-temperature superconductivity and strongly correlated systems, & topological quantum computing.
- Studying structure of scattering amplitudes of massive gauge bosons and gravitons in the Chern-Simons theories provides an important means for understanding the mechanism of **Topological Mass Generations** and for realizing Gauge-Gravity duality connection: **(Gravity) = (Gauge)<sup>2</sup>**

# 3d Topological Massive Gauge & Gravity Theories

➤ 3d Chern-Simons (CS) Topological Massive Gauge Theories:

$$\mathcal{L}_{\text{TMQED}} = -\frac{1}{4}F_{\mu\nu}^2 + \frac{1}{2}\tilde{m}\varepsilon^{\mu\nu\rho}A_{\mu}\partial_{\nu}A_{\rho},$$

$$\mathcal{L}_{\text{TMYM}} = -\frac{1}{2}\text{tr}F_{\mu\nu}^2 + \tilde{m}\varepsilon^{\mu\nu\rho}\text{tr}\left(A_{\mu}\partial_{\nu}A_{\rho} - \frac{i2g}{3}A_{\mu}A_{\nu}A_{\rho}\right),$$

➤ CS mass is geometrized and is related to CS Level:

$$n = 4\pi\tilde{m}/g^2 \in \mathbb{Z}$$

➤ 3d CS Topological Massive Gravity Theories:

$$S_{\text{TMG}} = -\frac{2}{\kappa^2}\int d^3x \left[ \sqrt{-g}R - \frac{1}{2\tilde{m}}\varepsilon^{\mu\nu\rho}\Gamma^{\alpha}_{\rho\beta}\left(\partial_{\mu}\Gamma^{\beta}_{\alpha\nu} + \frac{2}{3}\Gamma^{\beta}_{\mu\gamma}\Gamma^{\gamma}_{\nu\alpha}\right) \right]$$

E<sup>2</sup>

E<sup>3</sup>

# Topological Mass Generation

- Conversion of physical d.o.f between  $\mathbf{m} = \mathbf{0}$  and  $\mathbf{m} \neq \mathbf{0}$  :

$$A_P^a = \frac{1}{\sqrt{2}}(A_T^a + A_L^a) \quad \longrightarrow \quad A_T^a$$

- 2 orthogonal unphysical states:

$$A_X^a = \epsilon_X^\mu A_\mu^a = \frac{1}{\sqrt{2}}(A_T^a - A_L^a),$$
$$A_S^a = \epsilon_S^\mu A_\mu^a,$$

- The physical d.o.f is conserved under  $m \rightarrow 0$  limit:

$$1 = 1$$

- Then, we ask:

$$\mathcal{T}[A_P^{a_1}, \dots, A_P^{a_N}, \Phi] \stackrel{?}{=} \mathcal{T}[\tilde{A}_T^{a_1}, \dots, \tilde{A}_T^{a_N}, \Phi]$$

# Topological Mass Generation: TET

- Indeed we can derive a **new identity** to connect the Scattering Amplitudes:

$$\begin{aligned}\mathcal{T}[A_P^{a_1}, \dots, A_P^{a_N}, \Phi] &= \mathcal{T}[\tilde{A}_T^{a_1}, \dots, \tilde{A}_T^{a_N}, \Phi] + \mathcal{T}_v, \\ \mathcal{T}_v &= \sum_{j=1}^N \mathcal{T}[\tilde{v}^{a_1}, \dots, \tilde{v}^{a_j}, \tilde{A}_T^{a_{j+1}}, \dots, \tilde{A}_T^{a_N}, \Phi],\end{aligned}$$

- Under high energy expansion, we derive, **at S-matrix level**,  
**Topological Equivalence Theorem (TET)**:

$$\mathcal{T}[A_P^{a_1}, \dots, A_P^{a_N}, \Phi] = \mathcal{T}[\tilde{A}_T^{a_1}, \dots, \tilde{A}_T^{a_N}, \Phi] + \mathcal{O}\left(\frac{m}{E}\right),$$

# Topological Massive Gravity from Double-Copy

Hang, HJH, Shen, JHEP, 2110.05399

➤ E Cancellations:  $E^4 \rightarrow E^1$

Amplitude	$\times \bar{s}_0^2$	$\times \bar{s}_0^{3/2}$	$\times \bar{s}_0$
$\mathcal{M}_s$	$-\frac{99+28c_{2\theta}+c_{4\theta}}{1-c_{2\theta}}$	$-i14(15c_\theta+c_{3\theta})\csc\theta$	$-\frac{2(75+326c_{2\theta}+47c_{4\theta})}{1-c_{2\theta}}$
$\mathcal{M}_t$	$\frac{99+28c_{2\theta}+c_{4\theta}}{4(1-c_\theta)}$	$i(102+105c_\theta+70c_{2\theta}+7c_{3\theta}+4c_{4\theta})\csc\theta$	$\frac{75-107c_\theta+326c_{2\theta}+268c_{3\theta}+47c_{4\theta}+31c_{5\theta}}{1-c_{2\theta}}$
$\mathcal{M}_u$	$\frac{99+28c_{2\theta}+c_{4\theta}}{4(1+c_\theta)}$	$i(-102+105c_\theta-70c_{2\theta}+7c_{3\theta}-4c_{4\theta})\csc\theta$	$\frac{75+107c_\theta+326c_{2\theta}-268c_{3\theta}+47c_{4\theta}-31c_{5\theta}}{1-c_{2\theta}}$
Sum	0	0	0

➤ Under high energy expansion:

$$\mathcal{M}_0[4h_P] = -\frac{i\kappa^2}{2048} m s_0^{\frac{1}{2}} (494c_\theta + 19c_{3\theta} - c_{5\theta}) \csc^3\theta$$

➤ Energy Cancellations for 4 graviton scattering amplitudes:

$$E^4 \rightarrow E^0 \quad \longrightarrow \quad \mathcal{O}(E^{12}) \rightarrow \mathcal{O}(E^1), \quad (\text{for } \mathcal{E}_{h_p} = 4 \text{ in 3d TMG})$$

# Topological Mass Generation: Covariant TMG

- Introduce dilaton field via conformal factor:

$$g_{\mu\nu} = \bar{g}_{\mu\nu} e^{-\kappa\phi},$$

$$\mathcal{L}_{\text{TMG}}^\phi = \frac{-1}{\kappa^2} \left\{ \sqrt{-\bar{g}} e^{-\kappa\phi/2} \left[ R + \frac{\kappa^2}{2} \bar{g}_{\mu\nu} \partial^\mu \phi \partial^\nu \phi \right] - \frac{\varepsilon^{\mu\nu\rho}}{2\tilde{m}} \Gamma^\alpha_{\rho\beta} \left( \partial_\mu \Gamma^\beta_{\alpha\nu} + \frac{2}{3} \Gamma^\beta_{\mu\gamma} \Gamma^\gamma_{\nu\alpha} \right) \right\}$$

- Construct Gauge-fixing Terms:

$$\begin{aligned} \mathcal{L}_{\text{TMG}}^{\text{GF1}} &= \frac{1}{2\xi} (\mathcal{F}_{\text{GF1}}^\mu)^2, & \mathcal{F}_{\text{GF1}}^\mu &= \partial_\nu \bar{h}^{\mu\nu} - \frac{1}{2} \partial^\mu (\bar{h} - \xi\phi), \\ \mathcal{L}_{\text{TMG}}^{\text{GF2}} &= \frac{1}{2\zeta} (\mathcal{F}_{\text{GF2}}^\mu)^2, & \mathcal{F}_{\text{GF2}}^\mu &= \frac{1}{2} \partial^\mu (\bar{h} - \zeta\phi), \end{aligned}$$

- Landau Gauge Propagators:

$$\begin{aligned} \mathcal{D}_{\mu\nu\alpha\beta}^h(p) &= -\frac{mp^\rho}{2p^2(p^2+m^2)} \left( \varepsilon_{\rho\mu\alpha} \bar{\eta}_{\nu\beta} + \varepsilon_{\rho\mu\beta} \bar{\eta}_{\nu\alpha} + \varepsilon_{\rho\nu\alpha} \bar{\eta}_{\mu\beta} + \varepsilon_{\rho\nu\beta} \bar{\eta}_{\mu\alpha} \right) \\ &\quad + \frac{im^2}{p^2(p^2+m^2)} \left( \bar{\eta}_{\mu\alpha} \bar{\eta}_{\nu\beta} + \bar{\eta}_{\mu\beta} \bar{\eta}_{\nu\alpha} - \bar{\eta}_{\mu\nu} \bar{\eta}_{\alpha\beta} \right), \\ D^\phi(p) &= -\frac{i}{p^2}, \end{aligned}$$

1/p<sup>3</sup>

➤ Landau Gauge Propagators ( $\zeta = 0$  and  $\xi = 0$ ):

$$\mathcal{D}_{\mu\nu\alpha\beta}^h(p) = -\frac{mp^\rho}{2p^2(p^2+m^2)} \left( \varepsilon_{\rho\mu\alpha} \bar{\eta}_{\nu\beta} + \varepsilon_{\rho\mu\beta} \bar{\eta}_{\nu\alpha} + \varepsilon_{\rho\nu\alpha} \bar{\eta}_{\mu\beta} + \varepsilon_{\rho\nu\beta} \bar{\eta}_{\mu\alpha} \right) + \frac{im^2}{p^2(p^2+m^2)} \left( \bar{\eta}_{\mu\alpha} \bar{\eta}_{\nu\beta} + \bar{\eta}_{\mu\beta} \bar{\eta}_{\nu\alpha} - \bar{\eta}_{\mu\nu} \bar{\eta}_{\alpha\beta} \right),$$

$$D^\phi(p) = -\frac{i}{p^2},$$

$1/p^3$

➤ Unitary Gauge Propagator ( $\zeta = \infty$  and  $\xi = 0$ ):

$$D_{\mu\nu\alpha\beta}^{h(U)}(p) = \frac{-i(\bar{\eta}_{\mu\alpha} \bar{\eta}_{\nu\beta} + \bar{\eta}_{\mu\beta} \bar{\eta}_{\nu\alpha} - \bar{\eta}_{\mu\nu} \bar{\eta}_{\alpha\beta})}{p^2+m^2} + \frac{i\xi}{p^4} (\eta_{\mu\alpha} p_\nu p_\beta + \eta_{\mu\beta} p_\nu p_\alpha + \eta_{\nu\alpha} p_\mu p_\beta + \eta_{\nu\beta} p_\mu p_\alpha)$$

$$- \frac{mp^\rho}{2p^2(p^2+m^2)} \left( \varepsilon_{\rho\mu\alpha} \bar{\eta}_{\nu\beta} + \varepsilon_{\rho\mu\beta} \bar{\eta}_{\nu\alpha} + \varepsilon_{\rho\nu\alpha} \bar{\eta}_{\mu\beta} + \varepsilon_{\rho\nu\beta} \bar{\eta}_{\mu\alpha} \right)$$

$$+ \frac{i}{p^2} \left( \bar{\eta}_{\mu\alpha} \bar{\eta}_{\nu\beta} + \bar{\eta}_{\mu\beta} \bar{\eta}_{\nu\alpha} - 2\bar{\eta}_{\mu\nu} \bar{\eta}_{\alpha\beta} \right) - \frac{i2}{p^4} (\eta_{\mu\nu} p_\alpha p_\beta + \eta_{\alpha\beta} p_\mu p_\nu),$$

$1/p^2$

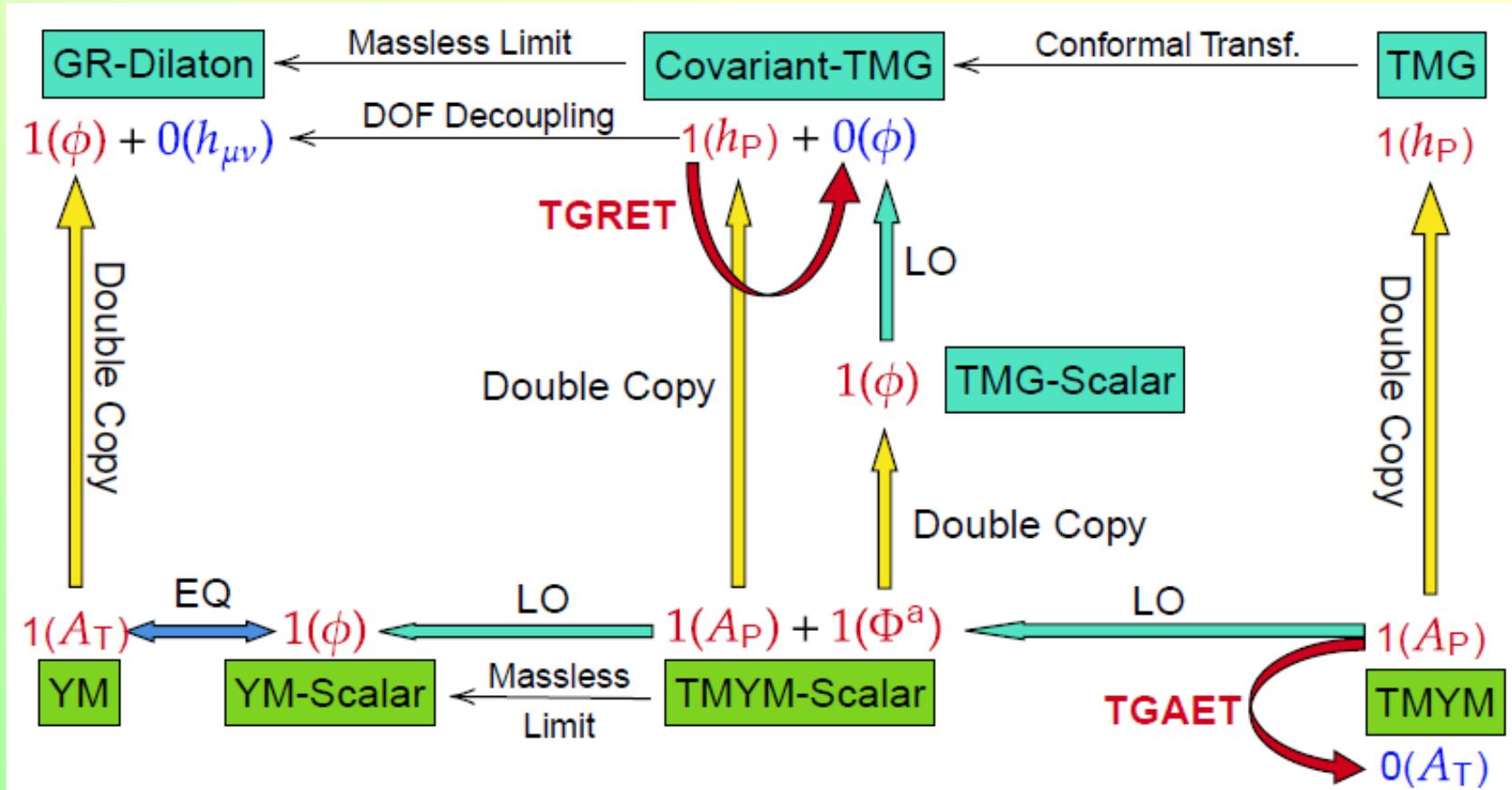
➤ Topological Gravity Equivalence Theorem (TGRET) :

$$\mathcal{M}[h_{\text{P}}(p_1), \dots, h_{\text{P}}(p_N), \Phi] = \left(\frac{1}{2}\right)^N C_{\text{mod}} \mathcal{M}[\phi(p_1), \dots, \phi(p_N), \Phi] + O\left(\frac{m}{E}\right),$$

➤ In comparison with TGAET:

$$\mathcal{T}[A_{\text{P}}^{a_1}(p_1), \dots, A_{\text{P}}^{a_N}(p_N), \Phi] = \left(\frac{1}{\sqrt{2}}\right)^N \mathcal{T}[A_{\text{T}}^{a_1}, \dots, A_{\text{T}}^{a_N}, \Phi] + O\left(\frac{m}{E}\right)$$

# Topological Mass Generation: Summary



# Generalized Power Counting for 3d Chern-Simons

- Generalized Energy Power Counting for 3d CS physical gauge boson scattering amplitudes:

$$D_E = (\mathcal{E}_{A_p} - \mathcal{E}_v) + (4 - \mathcal{E} - \bar{\mathcal{V}}_3) - L,$$

- Large E Cancellations in  $N$  physical  $A_p$  amplitudes by TET:

$$\Delta D_E = D_E[N A_p^a] - D_E[N A_T^a] = N$$

- E-Power Counting of  $N$  physical Graviton Amplitudes:

$$D_E[N h_P] = 2N + 2 + (\mathcal{V} - I_h + L) = 2N + 3$$

$$D_E^U[N h_P] = (2N + 2) + \mathcal{V} + L = (2N + 3) + I_h$$

$$D_{E(0)}[N \phi] = 3 - n = 3 - \frac{1}{2}N, \quad (N = 2n),$$

$$D_{E(0)}[N \phi] = 2 - n = \frac{1}{2}(5 - N), \quad (N = 2n + 1),$$

**N=4 Scattering Amplitudes:**

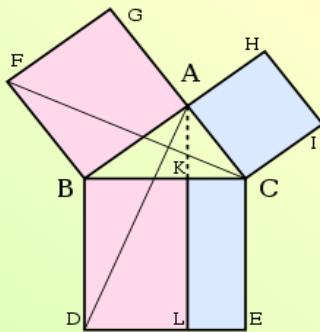
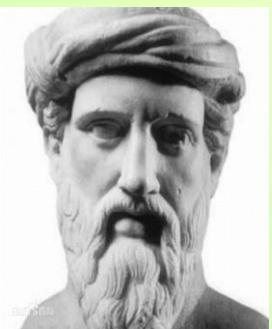
**$E^{12} \rightarrow E^1$  (Unitary Gauge)**

**$E^{11} \rightarrow E^1$  (Landau Gauge)**

$$D_{E(0)}[N h_P] - D_{E(0)}[N \phi] = \begin{cases} \frac{5}{2}N, & (\text{for } N = \text{even}), \\ \frac{1}{2}(5N + 1), & (\text{for } N = \text{odd}), \end{cases}$$

$$D_{E(0)}^U[N h_P] - D_{E(0)}[N \phi] = \begin{cases} \frac{7}{2}N - 3, & (\text{for } N = \text{even}), \\ \frac{1}{2}(7N - 5), & (\text{for } N = \text{odd}). \end{cases}$$

# Deep Relations: Importance of Equality & Square Law



- [EG] Pythagoras Theorem:  $a^2 + b^2 = c^2$  → Fermat Last Theorem:  $a^n + b^n \neq c^n$  ( $n > 2$ )
- [SR] Mass-Energy Equivalence:  $E = Mc^2$  and  $P_\mu^2 = M^2c^2$
- [GR] Inertial Mass vs Gravitational-Mass Equivalence:  $M_G = M_I$  →  $g = a$
- Gravity Force from Gauge Force: (valid for  $M = 0$  and  $M \neq 0$ )

→   
  $M = 0$   
(Conventional works)

$$(\text{Gravity}) = (\text{Gauge Force})^2$$

←   
  $M \neq 0$   
(This Talk)



→ **Much more can be explored !**

***Thank you***